A Measurement of the Electric Form Factor of the Neutron through $\vec{d}(\vec{e}, e'n)p$ at $Q^2 = 0.5$ (GeV/c)^2

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We report the first ever measurement using a solid polarized target of the neutron electric form factor G_E^n via $\vec{d}(\vec{e}, e'n)p$. G_E^n was determined from the beam–target asymmetry in the scattering of longitudinally polarized electrons from polarized deuterated ammonia ($^{15}ND_3$). The measurement was performed in Hall C at Thomas Jefferson National Accelerator Facility (TJNAF) in quasi free kinematics with the target polarization perpendicular to the momentum transfer. The electrons were detected in a magnetic spectrometer in coincidence with neutrons in a large solid angle segmented detector. We find $G_E^n = 0.04632 \pm 0.00616(\text{stat.}) \pm 0.00341(\text{syst.})$ at $Q^2 = 0.495$ (GeV/c)^2.

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Precise data on the neutron (and proton) form factors are important for understanding the non–perturbative mechanism responsible for confinement and are necessary in the interpretation of the electromagnetic properties of nuclei. The magnetic form factor of the neutron [1] has recently been measured with high precision. The neutron charge form factor G_E^n , in contrast, is only now yielding to intense efforts focused on its determination.

The major difficulty faced in a measurement of the neutron form factors is the lack of a free neutron target. The determination of G_E^n is further impeded by its small size. Advances in polarized electron sources, CW accelerators, polarimeters and polarized targets, now allow G_E^n to be extracted from experiments which exploit spin degrees of freedom. In particular the interference of the magnetic and electric scattering amplitudes is responsible for an asymmetry that can be measured in both polarized electron/polarized target experiments ($\vec{d}(\vec{e}, e'n)p$ [2] and $^3\vec{\mathrm{He}}(\vec{e}, e'n)p$ [3–5]) and in polarized electron recoil polarization measurements ($d(\vec{e}, e'\vec{n})p$ [6–8]).

For a vector polarized target of free neutrons, with the polarization P_n, in the scattering plane and perpendicular to the momentum transfer $\vec{q}$, G_E^n is related to the helicity asymmetry A_{en}^V [9] by

$$A_{en}^V = \frac{-2\sqrt{\tau(\tau+1)}\tan(\theta_e/2) G_E^n G_M^n}{(G_E^n)^2 + \tau[1 + 2(1+\tau)\tan^2(\theta_e/2)](G_M^n)^2} \quad (1)$$

where Q^2 is the four momentum transfer, $\tau = Q^2/4M_n^2$, and θ_e is the electron scattering angle. A_{en}^V is related to the experimental counts asymmetry $\epsilon = (L-R)/(L+R)$, where L and R are charge normalized counts for opposite beam polarizations P_e, by $A_{en}^V = \epsilon/(P_e P_n f)$, where f is the dilution factor due to scattering from materials other

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than polarized neutrons.

In practice one measures the helicity asymmetry from polarized deuterons in quasi elastic kinematics. The counts asymmetry for polarized electron polarized deuteron scattering can be written [10, 11] as:

$$\epsilon = f\frac{P_e A_e + P_e P_1^d A_{ed}^V + P_e P_2^d A_{ed}^T}{1 + P_1^d A_d^V + P_2^d A_d^T} \qquad (2)$$

where $P_{1(2)}^d$ is the target vector (tensor) polarization, and $A_e, A_d^V, A_d^T, A_{ed}^V$, and A_{ed}^T are the electron beam induced asymmetry, the vector and tensor deuteron target asymmetries, and the deuteron vector and tensor beam–target asymmetries, respectively. For experiments with the target polarization in the scattering plane which sample the neutron Fermi cone in an azimuthally symmetric way this reduces to

$$\epsilon = f\frac{P_e P_1^d A_{ed}^V}{1 + P_2^d A_d^T}. \qquad (3)$$

Appropriate averaging of data where the the sign of P_1^d is reversed reduces the sensitivity to A_e and A_{ed}^T if the setup does not have perfect azimuthal symmetry. For most practical targets P_2^d is small (3%) and the second term in the denominator may be neglected. Realistic calculations indicate that A_{ed}^V has a linear sensitivity to the magnitude of G_E^n for $d(e, e'n)$ at low recoil momentum [10, 11].

We present in this letter a measurement at $Q^2 = 0.495$ $(\mathrm{GeV/c})^2$ carried out at TJNAF, where the accelerator provides CW longitudinally polarized electrons. This experiment took place in Hall C with an incident electron energy of 2.725 GeV and a beam current of $\approx$ 100 nA incident on a dynamically polarized solid deuterated ammonia target. The scattered electrons were detected in coincidence with the knockout neutrons.

The polarized electron beam was produced by photoemission from a strained-layer semiconductor cathode illuminated by circularly polarized laser light [12] at the accelerator injector. The helicity of the beam was changed in a pair–wise pseudo-random sequence once per second to minimize sensitivity to instrumental drifts. The longitudinal polarization of the electrons was measured at regular intervals during the experiment with a Møller polarimeter [13] just upstream of the target. The average beam polarization for the data taking was $P_e = 0.776 \pm 0.002$ (stat.).

To prevent localized heating of the target material and to insure uniform irradiation, the beam was rastered over the face of the target such that the full face of the target was illuminated during each helicity state. The beam position was recorded by a secondary emission monitor [14] consisting of thin stainless steel strips in both the horizontal and vertical directions. It gave the transverse position of the interaction point of the electron and provided a calibration of the raster magnet system so that the beam position at the target could be used in the reconstruction.

The polarized target [15] included a permeable target cell filled with granules of $^{15}ND_3$ submerged in liquid He maintained at 1K by a high power He evaporation refrigerator. A 5 T magnetic field was provided by a superconducting coil arranged as a Helmholtz pair. The magnetic field was in the horizontal plane perpendicular to $< \vec{q} >$ at 151.6° with respect to the beam direction. The field orientation was measured in situ with a Hall probe to $\pm 0.1°$. A three magnet chicane compensated for the effects of the target magnetic field on the incident electrons. The field effects on the scattered electrons tilted the scattering plane by 4° with respect to the horizontal plane. The $^{15}ND_3$ had been previously irradiated at either Stanford or at the TJNAF Free Electron Laser with low energy electrons ($\simeq$ 35 MeV) to introduce a dilute collection of paramagnetic centers. The material was polarized by the dynamic nuclear polarization method. The polarization was measured continuously via NMR using a series LCR circuit and Q meter detector [16]. The average deuteron polarization throughout the experiment was $P_1^d = 0.21 \pm 0.01$.

The High Momentum Spectrometer (HMS) in its standard configuration was set at 15.7° to detect the scattered electrons. Modifications were made to the standard reconstruction algorithm of the HMS to account for the target field and the beam raster offset.

Knockout nucleons (p and n) were detected in an array of plastic scintillators. It consisted of 2 planes of thin (0.6 cm) veto paddles and 5 planes of 10 cm thick bars. The scintillator bars (160 cm long in the horizontal direction) had a phototube at each end to allow good position and timing resolution. The detector was positioned at 61.6° (along the direction of $< \vec{q} >$), providing a solid angle of $\approx$ 160 msr and was enclosed in a thick concrete walled hut open towards the target. The front shielding consisted of 16.7 mm of lead and 25 mm of CH_2 sheets. The time resolution was 450 ps (σ) as determined from the time of flight peak of photons (from π^0 decay) in the meantime spectrum. With the detector positioned 4.0 m from the target for nucleons with kinetic energy of 267 MeV it provided an energy resolution of 16.5 MeV. The neutron energy combined with the scattered electron energy allowed us to eliminate events associated with pion production. The neutron vertical position was determined by the segmentation of the detector (10 cm) while the horizontal position was determined from the time difference of the phototubes on the first bar hit along the n track. The measured horizontal resolution was $\approx$ 5 cm.

The electron-nucleon trigger was formed by a coincidence between the HMS electron and a hit in any one of the 5 bar planes. Neutrons were identified as events with no hits in the paddles along the track to the target, within a narrow time interval, and in a narrow range of invariant mass W around the quasi elastic peak ($\mid W - 938\mathrm{MeV} \mid < 50$ MeV). In addition, cuts on the horizontal position ($|y_{pos}| < 40$ cm) in the neutron detector and on the angle between $\vec{q}$ and the neutron momentum

($\theta_{nq} < 110$ mrad) were applied to optimize the dilution factor. The θ_{nq} cut served also to limit the recoil momentum p_r to values where the model dependence of A_{ed}^V has been shown to be small ($p_r^{max} \approx 85$ MeV/c) [10]. The protons were bent vertically in the target field by nearly 18° almost eliminating their overlap with the neutrons which further improved their rejection.

The experimental asymmetry was diluted by scattering from materials other than polarized deuterium nuclei. These include the nitrogen in $^{15}ND_3$, the liquid helium, the NMR coils, and target windows. A Monte Carlo (MC) simulation program was developed [17] to determine the dilution factor and to perform the detector averaging of the theoretical asymmetries. It was based on MCEEP [18] and included an HMS model, the neutron detector geometry and approximate efficiencies, the target magnetic field, the beam raster and radiative effects. Quasi elastic scattering from all the target materials was simulated in the MC. The normalization was fixed by data from carbon (which approximates nitrogen) and liquid helium. A comparison of the simulated distributions to experimental data is shown in Fig. 1 for four kinematic variables. The good agreement of the distributions indicates that quasi elastic scattering is the dominant process for events passing our selection criteria.

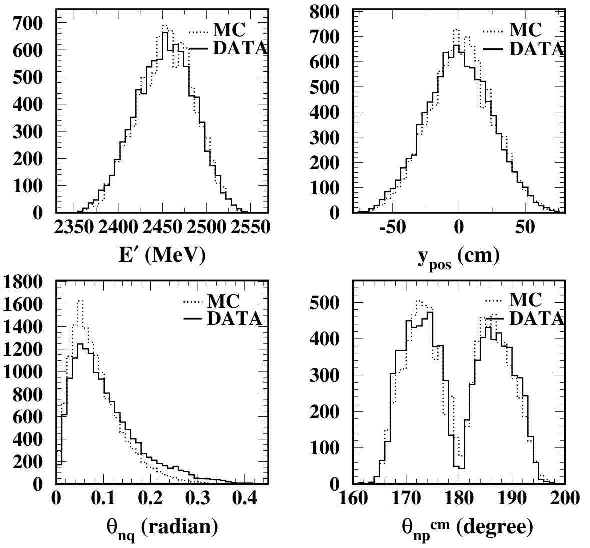

FIG. 1: A comparison between the data and the MC simulation for (e,e'n) from $^{15}ND_3$ for four kinematic variables: E' (scattered electron energy), $y_{\rm pos}$ (horizontal position in the neutron detector), θ_{nq} (the angle between $\vec{q}$ and the neutron in the lab), and θ_{np}^{cm} (the calculated angle between the proton and $\vec{q}$ in the n-p center of mass).

The accidental background under the meantime distribution was 4% and had no statistically significant asymmetry. The measured asymmetry was corrected for this dilution. A correction of 0.2% was made for the proton contamination. Charge exchange in deuterium was taken into account in the theoretical calculations of final state interactions (FSI). No correction was applied for charge exchange reactions with other target materials or the shielding (estimated to be 0.24%). The role of radiative effects on A_{ed}^V was estimated and found to be small. Corrections to the asymmetry for internal radiative effects of 2% and 0.5% for external effects were applied.

In order to extract G_E^n, the corrected experimental asymmetry was compared to the MC simulation which weights theoretical calculations of the asymmetry by the event distribution. The theoretical asymmetries have been calculated on a grid reflecting our experimental arrangement under different assumptions for the size of G_E^n. Asymmetry values between grid points were obtained by interpolation.

The theoretical A_{ed}^V values were calculated following [10, 11]. The calculations are based on a non-relativistic description of the n–p system in the deuteron, using the Bonn R-Space NN potential [19] for both the bound state and the description of FSI. The full calculations include meson exchange currents and isobar configurations as well as relativistic corrections. The dipole parameterization for G_M^n was assumed. It was verified that the acceptance averaged value of A_{ed}^V is linear in the size of G_M^n. Thus any (experimental) value could be incorporated easily if desired. The grid of asymmetries was calculated for 3 values of G_E^n. In each case the Q^2 variation of G_E^n was assumed to be given by the Galster parameterization [20](with $p = 5.6$) with the magnitude set by an overall scale parameter of 0.5, 1 or 1.5. The narrow acceptance in Q^2, $0.4 < Q^2 < 0.6$ (GeV/c)^2, makes the extracted value of G_E^n insensitive to the assumed Q^2 dependence.

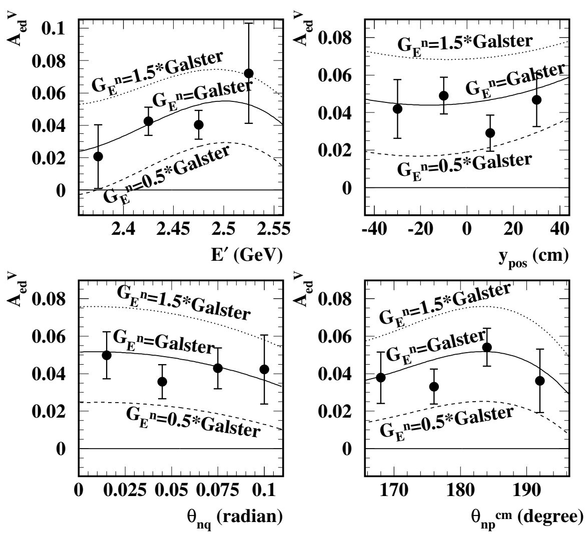

FIG. 2: Comparison between the measured asymmetries and calculated values of A_{ed}^V for three scaled parameterizations of Galster shown against four kinematic variables (See Fig. 1).

The MC simulation averaged the asymmetry over all kinematic variables except the one under investigation. The comparisons are shown in Fig. 2. To determine the best value of the scale parameter, it was fit as a free parameter to the data. The resulting value for G_E^n at $Q^2 = 0.495$ (GeV/c)^2 is $G_E^n = 0.04632 \pm 0.00616(\text{stat.}) \pm 0.00341(\text{syst.})$.

The major sources of systematic errors are: $\delta P_1^d/P_1^d = 5.8\%$, and the uncertainty in determining the average dilution factor is 3.9%. Cut dependencies give a 2.4% systematic error. Errors associated with the determination of various kinematic quantities contribute another 2.2%. The determination of P_e contributes 1%. Finally, the error in the value of G_M^n was taken to be 1.7% as derived from a recent fit to world data [21]. The quadratic sum of all the contributions gives a total systematic error $\delta G_E^n/G_E^n = 7.4\%$.

Our measurement is compared to G_E^n measurements from other polarized experiments in Fig. 3. The solid line and band represents our fit of the data set shown with the form

$$G_E^n(Q^2) = -\mu_n \frac{a\tau}{1+p\tau} G_D(Q^2)$$

with $G_D(Q^2) = 1/(1+Q^2/\Lambda^2)^2$ and $\Lambda^2 = 0.71(\text{GeV/c})^2$. The paramter a is determined by a measurement of $\frac{dG_E^n}{dQ^2}$ at $Q^2 = 0$ [22]. We find, for $a = 0.895 \pm 0.039$ and $p = 4.36 \pm 1.11$ for the six data points. The figure shows also the results of recent lattice QCD calculations [23]. In these calculations G_E^n is quite sensitive to the disconnected insertions which account for the sea-quarks. The magnitude of these sea contributions to the various nucleon form factors is almost constant so they are relatively much more important for G_E^n. Thus G_E^n may provide a valuable testing ground for lattice calculations of other sea sensitive quantities such as the strangeness electric and magnetic form factors.

The size of reaction dynamic effects beyond the plane wave Born approximation (PWBA) was determined by repeating the same extraction procedure using PWBA calculations. The result for G_E^n was found to be 13% smaller than when it was extracted from A_{ed}^V using the full calculation. The bulk of the difference is due to FSI.

In conclusion we present the results of a new measurement of the neutron electric form factor at $Q^2 = 0.495$ (GeV/c)^2, the highest momentum transfer to date in polarized scattering using a deuteron target. This measurement sets a new constraint on the parameterizations of G_E^n and, more importantly, on theoretical models which describe it. In addition it will contribute to the extraction of strange quark form factors from parity violating (PV) elastic scattering from protons [25, 26] where errors on previous measurements of G_E^n are the largest contributor to the theoretical PV asymmetry, A_{th}. The ongoing effort to measure G_E^n will considerably extend our understanding of nucleon structure.

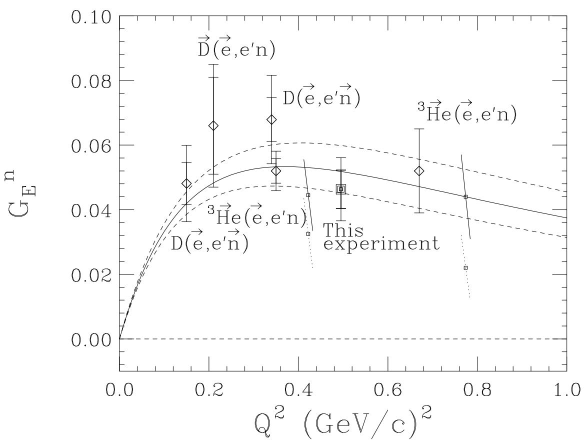

FIG. 3: Comparison of present experiment (filled square) with data from recent polarized scattering measurements, left to right [7], [2], [8], [4], [3]. FSI corrections have been made for all the data except [3] where FSI are thought to be modest. Data from Ref. [4] have been corrected for FSI [24]. The solid line and band represents our fit of this data set (see text). The slanted lines are lattice QCD calculations with (solid lines) and without (dotted lines) the disconnected insertions which account for sea-quark effects [23].

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