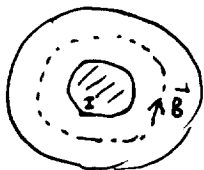


From the cylindrical symmetry, we know that the mag. field will be tangent to a circle centered on the axis of the cylinders and will depend only on the distance from the axis.

We use Ampere's law:  $\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{enclosed}}$



$$\Rightarrow B \cdot 2\pi r = \mu_0 I$$

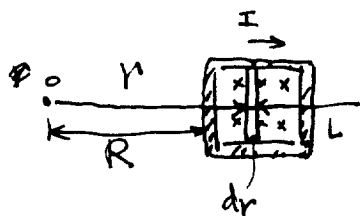
$$\Rightarrow B = \frac{4\pi \times 10^{-7} \times 10}{2\pi \times r}$$

Thus at the midway where  $r = \frac{0.1 + 0.5}{2} = 0.3 \text{ cm}$

$$B = \frac{4\pi \times 10^{-7} \times 10}{2\pi \times 0.3 \times 10^{-2}} = 6.7 \times 10^{-4} \text{ T in circular direction.}$$

#31. The magnetic field inside the toroidal solenoid is circular and varies with the distance from the center of the torus  $r$ :

$$B = \frac{\mu_0 N I}{2\pi r} \quad (\text{from Ampere's law})$$



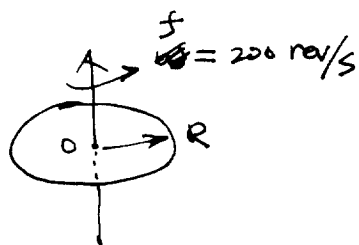
We find flux by integration.

$$\Phi_B = \iint \vec{B} \cdot d\vec{A} = \int_R^{R+L} B(r) \cdot L dr$$

$$= \int_R^{R+L} \frac{\mu_0 N I}{2\pi r} \cdot L dr = \frac{\mu_0 N I L}{2\pi} \int_R^{R+L} \frac{dr}{r}$$

$$\Rightarrow \Phi_B = \frac{\mu_0 N I L}{2\pi} \ln\left(\frac{R+L}{R}\right)$$

#38)



the rotating ring is equivalent to a circular current, which has a magnitude:

$$I = Q/T = Qf$$

thus the magnetic field at the center of the ring is:

$$B = \mu_0 I / 2R = \mu_0 Qf / 2R = 4\pi \times 10^{-7} \times 2 \times 10^{-5} \times 200 / 2 \times 1.5 \times 10^{-2} = 1.7 \times 10^{-7} \text{ T}$$

along the axis.

#43



Because the point P is along the line of the two straight segments of the wire, there is no magnetic field from these segments.

The magnetic field at the point P is the sum of

$$\text{two fields: } \vec{B} = \vec{B}_{\text{inner semi-circle}} + \vec{B}_{\text{outer semi-circle}}$$

Each field is half that of a circular loop, with the field of the inner semicircle into the page and that of the outer semicircle out of the page, so:

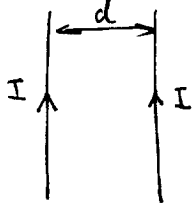
$$B = \frac{1}{2} \frac{\mu_0 I}{2R_{\text{inner}}} - \frac{1}{2} \frac{\mu_0 I}{2R_{\text{outer}}} = \frac{\mu_0 I}{4} \cdot \frac{1}{R_{\text{inner}}} - \frac{\mu_0 I}{4} \frac{1}{R_{\text{outer}}}$$

$$= \pi \times 10^{-7} \times 12 \left( \frac{1}{0.05} - \frac{1}{0.08} \right) = 2.8 \times 10^{-5} \text{ T into the page.}$$

#63)

a) For the attractive magnetic force per unit length,

we have  $F_B/L = \frac{I \mu_0 I}{2\pi d} = \frac{2 \times 10^{-7} \times 1^2}{1 \times 10^{-2}} = 2.0 \times 10^{-5} \text{ N/m}$



the linear charge density is:

$$\lambda = 10^{21} \times 100 \times 1.6 \times 10^{-19} = 1.6 \times 10^4 \text{ C/m}$$

The electric field  $E$  at a distance  $d$  from an infinite line of charge is  $E = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{d}$  (22-31)

For the repulsive electric force per unit length, we have

$$F_E/L = \lambda E = \frac{1}{2\pi\epsilon_0} \frac{\lambda^2}{d} = 2 \times \frac{9 \times 10^9 \times 1.6 \times 10^4}{1 \times 10^{-2}} = 4.6 \times 10^{20} \text{ N/m}$$

the ratio of the forces is:

$$\frac{F_B}{F_E} = 4.3 \times 10^{-26}$$

b) For the forces to be equal, we have:

$$\frac{\lambda^2}{\epsilon_0} = \mu_0 I^2 \quad \text{or} \quad \lambda^2 = \epsilon_0 \mu_0 I^2$$

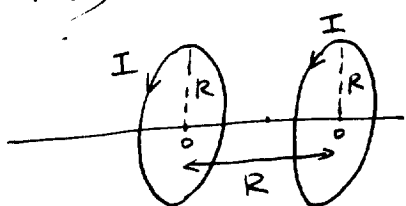
since  $\epsilon_0 \mu_0 = 1/c^2$  and  $\lambda = ne$ , we get:

$$n = I/ce = \frac{1}{(3 \times 10^8) \times (1.6 \times 10^{-19})} = 2.1 \times 10^{10} \text{ electrons/m} = 2.1 \times 10^8 \text{ /cm}$$

c) For the fraction that is the excess, we have:

$$f = \frac{2.1 \times 10^8}{10^{21}} = 2.1 \times 10^{-13}$$

#65)

(a) We choose  $x=0$  at the left coil.

The magnetic fields on the axis from the coils are in the same direction, so we find the total magnetic field from:

$$B(x) = \frac{\mu_0 I}{2} \left\{ \frac{R^2}{(R^2 + x^2)^{3/2}} + \frac{R^2}{[R^2 + (R-x)^2]^{3/2}} \right\}$$

$$= \frac{\mu_0 I}{2R} \left\{ \frac{1}{[1 + (\frac{x}{R})^2]^{3/2}} + \frac{1}{[2 - 2(\frac{x}{R}) + (\frac{x}{R})^2]^{3/2}} \right\}$$

at  $x=0$ ,  $B(0) = \frac{\mu_0 I}{2R} \left( 1 + \frac{1}{2^{3/2}} \right) = 0.677 \mu_0 I/R$

at  $x=R/4$ :  $B(R/4) = \frac{\mu_0 I}{2R} \left( \frac{1}{[1 + (\frac{1}{4})^2]^{3/2}} + \frac{1}{[2 - 2 \times \frac{1}{4} + (\frac{1}{4})^2]^{3/2}} \right) = 0.713 \mu_0 I/R$

at  $x=R/2$ :  $B(R/2) = \frac{\mu_0 I}{2R} \left( \frac{1}{[1 + (\frac{1}{2})^2]^{3/2}} + \frac{1}{[2 - 2 \times \frac{1}{2} + (\frac{1}{2})^2]^{3/2}} \right) = 0.716 \mu_0 I/R$

(b)  $\frac{dB}{dx} = -\frac{3\mu_0 I}{2R^2} \left\{ \frac{x/R}{[1 + (\frac{x}{R})^2]^{5/2}} + \frac{x/R - 1}{[2 - 2(\frac{x}{R}) + (\frac{x}{R})^2]^{5/2}} \right\}$

$$\frac{d^2B}{dx^2} = -\frac{3\mu_0 I}{2R^3} \left\{ \frac{1}{[1 + (\frac{x}{R})^2]^{5/2}} - \frac{5(x/R)^2}{[1 + (\frac{x}{R})^2]^{7/2}} + \frac{1}{[2 - 2(\frac{x}{R}) + (\frac{x}{R})^2]^{5/2}} - \frac{5[x/R - 1]^2}{[2 - 2(\frac{x}{R}) + (\frac{x}{R})^2]^{7/2}} \right\}$$

$\therefore$  At  $x=R/2$ :

$$\frac{dB}{dx} = -\frac{3\mu_0 I}{2R^2} \left\{ \frac{1/2}{[1 + (\frac{1}{2})^2]^{5/2}} + \frac{1/2 - 1}{[2 - 2(\frac{1}{2}) + (\frac{1}{2})^2]^{5/2}} \right\} = 0$$

$$\frac{d^2B}{dx^2} = -\frac{3\mu_0 I}{2R^3} \left\{ \frac{1}{[1 + (\frac{1}{2})^2]^{5/2}} - \frac{5(\frac{1}{2})^2}{[1 + (\frac{1}{2})^2]^{7/2}} + \frac{1}{[2 - 2(\frac{1}{2}) + (\frac{1}{2})^2]^{5/2}} - \frac{5[1/2 - 1]^2}{[2 - 2(\frac{1}{2}) + (\frac{1}{2})^2]^{7/2}} \right\}$$

$$= 0$$