

Lecture 3 - Identical Particles

Start with Ch 5 - Identical Particles

Begin by consider quantum system of two particles

Particle 1 position \vec{r}_1
 " 2 " \vec{r}_2

Describe with two particle wavefunction

$$\psi(\vec{r}_1, \vec{r}_2, t)$$

Satisfies Schrodinger equation

$$i\hbar \frac{\partial \psi}{\partial t} = H\psi$$

$$\text{with } H = -\frac{\hbar^2}{2m} \nabla_1^2 - \frac{\hbar^2}{2m} \nabla_2^2 + V(\vec{r}_1, \vec{r}_2, t)$$

$$\nabla_i^2 = \frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} + \frac{\partial^2}{\partial z_i^2}, \text{ sim for } \nabla_2$$

Normalization

$$\int |\psi|^2 d^3r_1 d^3r_2 = 1$$

Interpret $|\psi(\vec{r}_1, \vec{r}_2)|^2$ = prob density to find
 particle 1 at \vec{r}_1
 particle 2 at \vec{r}_2

[In general, have many body wave function
 $\psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N, t)$... same idea]

If H doesn't depend on time, can separate variables:

$$\psi(\vec{r}_1, \vec{r}_2, t) = \psi(\vec{r}_1, \vec{r}_2) e^{-iEt/\hbar}$$

with $-\frac{\hbar^2}{2m} \nabla_1^2 \psi - \frac{\hbar^2}{2m} \nabla_2^2 \psi + V\psi = E\psi$

* very hard to solve!

- If $V(\vec{r}_1, \vec{r}_2) = V(\vec{r}_1) + V(\vec{r}_2)$
particles are independent

can separate further: $\psi(\vec{r}_1, \vec{r}_2) = \psi_1(\vec{r}_1)\psi_2(\vec{r}_2)$

with $-\frac{\hbar^2}{2m} \nabla_1^2 \psi_1 + V(\vec{r}_1) \psi_1 = E_1 \psi_1$

$-\frac{\hbar^2}{2m} \nabla_2^2 \psi_2 + V(\vec{r}_2) \psi_2 = E_2 \psi_2$

$$E = E_1 + E_2$$

Just two copies of single-particle problem

- More interesting when particles interact:

Typically $V(\vec{r}_1, \vec{r}_2) = V(\vec{r}_1 - \vec{r}_2)$

(no external potential)

Then can't separate $\vec{r}_1 + \vec{r}_2$

Can separate different variables:

center-of-mass coordinates.

[Problem 5.1]

Define $\vec{r} = \vec{r}_1 - \vec{r}_2$ relative coordinate

$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}$ COM coordinate

Convert Schr. eqn. to these coordinates

$$\vec{r}_2 = \vec{r}_1 - \vec{r}$$

$$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 (\vec{r}_1 - \vec{r})}{m_1 + m_2} = \vec{r}_1 - \frac{m_2}{m_1 + m_2} \vec{r}$$

So $\boxed{\vec{r}_1 = \vec{R} + \frac{m_2}{m_1 + m_2} \vec{r}}$

$$\vec{r}_2 = \vec{r}_1 - \vec{r} = \vec{R} + \left(\frac{m_2}{m_1 + m_2} - 1 \right) \vec{r}$$

$\boxed{\vec{r}_2 = \vec{R} - \frac{m_1}{m_1 + m_2} \vec{r}}$

Define $M = m_1 + m_2$
= total mass

Need to convert derivatives $\vec{\nabla}_1 + \vec{\nabla}_2$

Have $\frac{\partial}{\partial x_1} = \frac{\partial X}{\partial x_1} \frac{\partial}{\partial X} + \frac{\partial x}{\partial x_1} \frac{\partial}{\partial x}$

$$X = \frac{m_1 x_1 + m_2 x_2}{M} \quad x = x_1 - x_2$$

$$\text{So } \frac{\partial}{\partial x_1} = \frac{m_1}{M} \frac{\partial}{\partial X} + \frac{\partial}{\partial x}$$

Similarly for other coordinates, set

$$\vec{\nabla}_1 = \frac{m_1}{M} \vec{\nabla}_R + \vec{\nabla}_r$$

$$\vec{\nabla}_2 = \frac{m_2}{M} \vec{\nabla}_R - \vec{\nabla}_r$$

$$\text{So } \vec{\nabla}_1^2 = \frac{m_1^2}{M^2} \vec{\nabla}_R^2 + \vec{\nabla}_r^2 + 2 \frac{m_1}{M} \vec{\nabla}_R \cdot \vec{\nabla}_r$$

$$\vec{\nabla}_2^2 = \frac{m_2^2}{M^2} \vec{\nabla}_R^2 + \vec{\nabla}_r^2 - 2 \frac{m_2}{M} \vec{\nabla}_R \cdot \vec{\nabla}_r$$

Put into Schr. Eqn:

$$\left[-\frac{\hbar^2}{2m_1} \nabla_1^2 - \frac{\hbar^2}{2m_2} \nabla_2^2 + V(\vec{r}) \right] \psi = E \psi$$

$$\left[-\frac{\hbar^2}{2} \frac{m_1}{M^2} \nabla_R^2 - \frac{\hbar^2}{2m_1} \nabla_r^2 - \frac{\hbar^2}{m} \vec{\nabla}_R \cdot \vec{\nabla}_r \right]$$

$$-\frac{\hbar^2}{2} \frac{m_2}{M^2} \nabla_R^2 - \frac{\hbar^2}{2m_2} \nabla_r^2 + \frac{\hbar^2}{m} \vec{\nabla}_R \cdot \vec{\nabla}_r \] = (E - V) \psi$$

or

$$\left[-\frac{\hbar^2}{2m} \nabla_R^2 - \frac{\hbar^2}{2} \left(\frac{1}{m_1} + \frac{1}{m_2} \right) \nabla_r^2 + V(\vec{r}) \right] \psi = E \psi$$

$$\text{Define } \frac{1}{\mu} = \frac{1}{m_1} + \frac{1}{m_2}$$

μ = reduced mass

Then separate:

$$-\frac{\hbar^2}{2m} \nabla_R^2 \psi = E_R \psi \rightarrow \text{free particle equation}$$

$$-\frac{\hbar^2}{2m} \nabla_r^2 \psi + V(\vec{r}) \psi = E_r \psi \rightarrow \text{single particle in potential}$$

$$\text{Total } E = E_r + E_R$$

Usually don't even think about \vec{R} ; $k_{\vec{R}}$
Know solutions are plane waves $\sim e^{i k_{\vec{R}} \cdot \vec{r}}$

You already know how to solve all the single particle problems that have exact solutions.

So you can apply that directly to two-body systems.

Note: doesn't work if there's an external potential $V(\vec{R})$!

Note, in general, particles will also have spin

$$\Psi(\vec{r}_1, \vec{r}_2, t) \rightarrow \Psi(\vec{r}_1, \vec{s}_1, \vec{r}_2, \vec{s}_2, t)$$

Often (not always) the case that spatial
+ spin parts separate

$$\rightarrow \Psi(\vec{r}_1, \vec{r}_2, t) \propto (\vec{s}_1, \vec{s}_2, t)$$

Then above discussion applies to spatial part

Use techniques from Ch 4 to handle spin part
(Know how to combine two spins)

Exchange Symmetry

Now suppose our particles are the same type
two electrons, for instance.

Then particles are indistinguishable: can't tell
one from the other

Must have $H(\vec{r}_1, \vec{r}_2) = H(\vec{r}_2, \vec{r}_1)$
symmetric under particle exchange

Use symmetry to define a new operator that
commutes with H :

P = exchange operator

$$P \Psi(\vec{r}_1, \vec{r}_2) = \Psi(\vec{r}_2, \vec{r}_1)$$

To see that P commutes with H :

$$\begin{aligned} PH^2 &= P[H(\vec{r}_1, \vec{r}_2) 2t(\vec{r}_1, \vec{r}_2)] = H(\vec{r}_2, \vec{r}_1) 2t(\vec{r}_2, \vec{r}_1) \\ &= H(\vec{r}_1, \vec{r}_2) 2t(\vec{r}_1, \vec{r}_2) \\ &= HP^2 \end{aligned}$$

$$so [P, H] = 0$$

So we can find eigenstates of H that are also eigenstates of P

What are eigenvalues of P ?

$$\text{Have } P^2 2t(\vec{r}_1, \vec{r}_2) = P 2t(\vec{r}_1, \vec{r}_2) = 2t(\vec{r}_1, \vec{r}_2)$$

$$so P^2 = 1$$

Suppose ϕ is an eigenstate of P

$$\text{Then } P\phi = \lambda\phi$$

$$P^2\phi = \lambda^2\phi = \phi$$

$$so \lambda^2 = 1 \Rightarrow \lambda = \pm 1$$

Two eigenstates: $\lambda = +1 \Rightarrow$ even exchange symmetry

$\lambda = -1 \Rightarrow$ odd exchange symmetry

Eigenstates have form

$$2t(\vec{r}_1, \vec{r}_2) = A[2t_a(\vec{r}_1) 2t_b(\vec{r}_2) \pm 2t_a(\vec{r}_2) 2t_b(\vec{r}_1)]$$

+ for even, - for odd