

Lecture 12

Two-level system: To understand perturbation theory

Say just two states: Basis $|4_a^0\rangle, |4_b^0\rangle$

$$\text{Hamiltonian } H = H^0 + H'$$

$$H^0 |4_a^0\rangle = E_a^0 |4_a^0\rangle$$

$$H^0 |4_b^0\rangle = E_b^0 |4_b^0\rangle$$

$$\text{Write general state } |\psi\rangle = \alpha |4_a^0\rangle + \beta |4_b^0\rangle$$

Convenient to write as vector $\begin{bmatrix} \alpha \\ \beta \end{bmatrix}$

Can write H^0 as matrix:

$$H^0 |\psi\rangle = \alpha E_a^0 |4_a^0\rangle + \beta E_b^0 |4_b^0\rangle$$

$$\text{or } H^0 = \begin{bmatrix} E_a^0 & 0 \\ 0 & E_b^0 \end{bmatrix}$$

What about H' ?

$$H' |\psi\rangle = \alpha H' |4_a^0\rangle + \beta H' |4_b^0\rangle$$

$$\text{So } \langle 4_a^0 | H' | \psi \rangle = \alpha \langle 4_a^0 | H' | 4_a^0 \rangle + \beta \langle 4_a^0 | H' | 4_b^0 \rangle$$

$$\text{and } \langle 4_b^0 | H' | \psi \rangle = \alpha \langle 4_b^0 | H' | 4_a^0 \rangle + \beta \langle 4_b^0 | H' | 4_b^0 \rangle$$

Integrals we can evaluate

Define $w_{aa} = \langle 24^\circ | H' | 24^\circ \rangle$ $w_{as} = \langle 24^\circ | H' | 24_s^\circ \rangle$

$$w_{ba} = \langle 24_s^\circ | H' | 24^\circ \rangle \quad w_{bb} = \langle 24_s^\circ | H' | 24_s^\circ \rangle \\ = w_{ab}^*$$

Then can write H' as matrix $\begin{bmatrix} w_{aa} & w_{as} \\ w_{ba} & w_{bb} \end{bmatrix}$

Schro. Eqn. $H|4\rangle = E|4\rangle$ becomes

$$\begin{bmatrix} E_a^0 + w_{aa} & w_{as} \\ w_{ba} & E_b^0 + w_{bb} \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix} = E \begin{bmatrix} \alpha \\ \beta \end{bmatrix}$$

Want to find eigenvalues E

Write $\begin{bmatrix} A & \\ \underbrace{\begin{bmatrix} E_a^0 + w_{aa} - E & w_{as} \\ w_{ba} & E_b^0 + w_{bb} - E \end{bmatrix}}_{B} \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix} = 0$

Determinant of matrix = 0

$$(A-E)(B-E) - |w_{as}|^2 = 0$$

$$E^2 - E(A+B) + AB - |w_{as}|^2 = 0$$

Quadratic:

$$E = \frac{1}{2} [A+B \pm \sqrt{(A+B)^2 - 4AB + 4|w_{as}|^2}]$$

$$\boxed{E = \frac{1}{2} [A+B \pm \sqrt{(A-B)^2 + 4|w_{as}|^2}]}$$

If $|W_{ab}| \ll |A-B|$, expand square root

$$E = \frac{1}{2} \left[A+B \pm (A-B) \left(1 + 2 \frac{|W_{ab}|^2}{(A-B)^2} \right) \right]$$

$$E_+ = A + \frac{|W_{ab}|^2}{A-B}$$

$$E_- = B - \frac{|W_{ab}|^2}{A-B}$$

or

$$E_a = E_a^0 + W_{aa} + \frac{|W_{ab}|^2}{E_a^0 - E_b^0} \quad \text{drop } W\text{'s to this order}$$

$$E_b = E_b^0 + W_{bb} + \frac{|W_{ab}|^2}{E_b^0 - E_a^0}$$

See this agrees with our perturbation theory results

$$E_a^1 = \langle \psi_a^0 | H' | \psi_a^0 \rangle = W_{aa}$$

$$E_a^2 = \frac{\langle \psi_a^0 | H' | \psi_b^0 \rangle \langle \psi_b^0 | H' | \psi_a^0 \rangle}{E_a^0 - E_b^0} = \frac{|W_{ab}|^2}{E_a^0 - E_b^0}$$

Funny behavior where 2nd-order term pushes energies apart:
because off-diagonal term in a matrix
pushes eigenvalues apart - not quite as mysterious

We use perturbation theory just because it's hard
to diagonalize a matrix bigger than 2x2

Now, what happens if $E_a^0 = E_b^0 \equiv E^0$?

Then
$$E = \frac{1}{2} \left[2E^0 + W_{aa} + W_{bb} \pm \sqrt{(W_{aa}-W_{bb})^2 + 4|W_{ab}|^2} \right]$$

$$= E^0 + E'$$

First order correction

$$E' = \frac{1}{2} \left[\omega_{aa} + \omega_{bb} \pm \sqrt{(\omega_{aa} - \omega_{bb})^2 + 4\omega_{ab}^2} \right]$$

can't simplify in general

However, if $\omega_{ab} = 0$, then $E' = \omega_{aa}$ or ω_{bb}
same as normal nondegenerate result

→ Good to pick $|E_a^0\rangle$ and $|E_b^0\rangle$ so that
 $\omega_{ab} = 0$

Since $E_a^0 = E_b^0$, can use any linear combination
of states for basis.

Trick: find some other operator Q that commutes
with H^0 and H'

(Parity, angular momentum, linear momentum
are good things to try)

Then use eigenstates of Q as basis:

since $[Q, H'] = 0$, guaranteed that $\omega_{ab} = 0$

If you can't find such an operator Q , just
use general formula

Note, sometimes have $0 < |E_a^0 - E_b^0| \approx \omega_{ab}$

Then must use general formula for E .
Really, perturbation theory doesn't apply