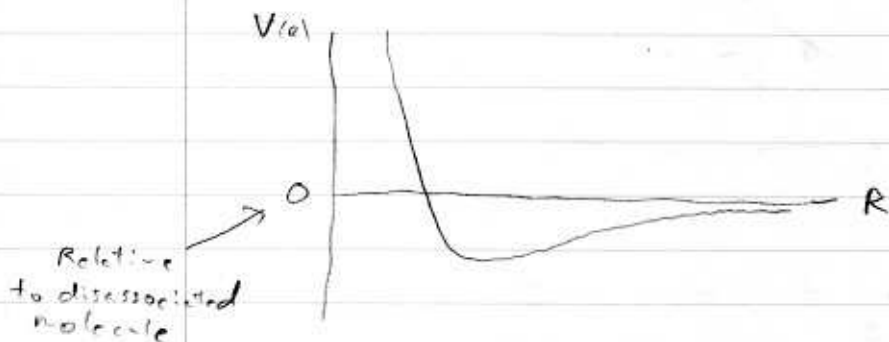


Lecture 24

Last time, applied variational method to H_2^+

Got estimates for bond length, binding energy

Also got estimate for molecular potential $V(R)$



Standard trick: use this as potential for nuclear motion:

$$H \rightarrow -\frac{\hbar^2}{2\mu} \nabla^2 \psi(R) + V(R)\psi(R) \quad \text{See Prob 7.10}$$

$\mu = \frac{M}{2}$, $M = \text{nuclear mass}$

Generally attack chemistry problems this way:

Step 1: Solve for electron energy at fixed nuclear positions

Step 2: Use e^- energy as potential for nuclear motion

Only works because electrons are so much lighter = faster
electron state adjusts adiabatically as nuclei move

Called Born-Oppenheimer approximation

Example of where this applies: collision between two atoms

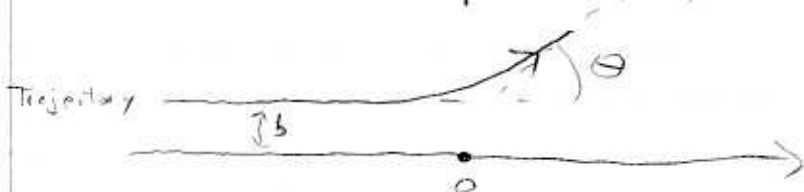
atoms interact via potential $V(r)$: typically complicated

Study how to approach this problem: Chapter 11

Scattering

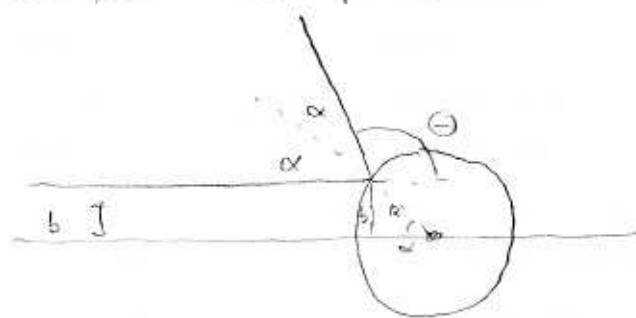
Start with a little classical scattering theory

Assume incident particle moving along z
scatters from spherically symmetric potential $V(r)$



Basic classical problem: determine scattering angle θ
as function of impact parameter b

Example: Hard sphere scattering:



Assume elastic collision

$$\text{See } b = R \sin \alpha$$

$$\text{and } 2\alpha + \theta = \pi$$
$$\alpha = \frac{\pi - \theta}{2}$$

$$\text{So } b = R \sin\left(\frac{\pi - \theta}{2}\right) = R \cos \frac{\theta}{2}$$

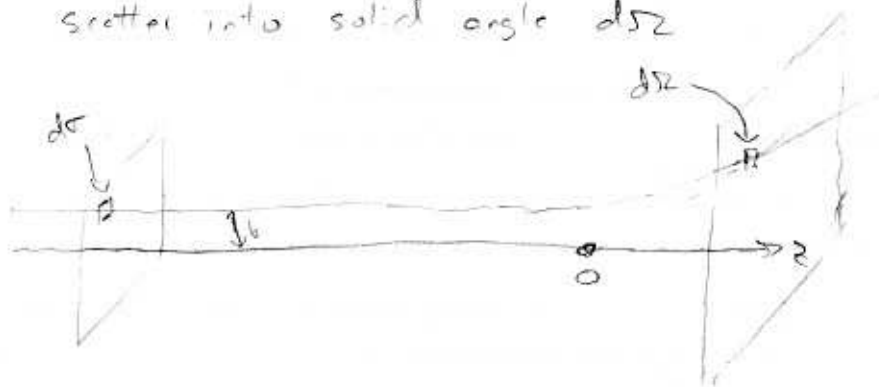
$$\text{and } \theta = \begin{cases} 2 \cos^{-1} \frac{b}{R} & (b \leq R) \\ 0 & (b > R) \end{cases}$$

Apply this to problem where we're sending
a beam of particles at scatterer.

Then we have a whole range of impact parameters

and generalize to 3D:

Particles incident within area $d\sigma$
 scatter into solid angle $d\Omega$



$$d\Omega = 2\pi \sin\theta d\theta$$

Of course range of solid angles is proportional to
 size of incident area

$$d\sigma = D(\theta) d\Omega$$

$D(\theta)$ = proportionality constant
 (only depends on θ for
 spherically symmetric v)

Can write $d\sigma = 2\pi b db = D(\theta) 2\pi \sin\theta d\theta$

$$\text{So } D(\theta) = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

For hard sphere, $b = R \cos \frac{\theta}{2}$

$$\frac{db}{d\theta} = -\frac{1}{2} R \sin \frac{\theta}{2}$$

$$D(\theta) = \frac{b}{\sin\theta} \cdot \frac{R}{2} \sin \frac{\theta}{2} = \frac{R^2}{2} \frac{\cos \frac{\theta}{2} \sin \frac{\theta}{2}}{\sin\theta}$$

But $\sin\theta = 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2}$

$$\text{So } D(\theta) = \frac{R^2}{4}$$

Note $D(\theta) = \frac{d\sigma}{d\Omega}$, units of area

Call D = differential cross section
after just write $d\sigma/d\Omega$

If we have a beam of particles with

$$L = \frac{\# \text{ of particles}}{\text{area} \cdot \text{time}} \equiv \text{luminosity or intensity}$$

Then scatter $dN = L D(\theta) d\Omega$ into solid angle $d\Omega$
per unit time

Makes D pretty easy to measure.

Conventional to define total cross section

$$\sigma = \int D(\theta) d\Omega \approx \text{effective area of target}$$

For hard sphere, $D(\theta) = \frac{R^2}{4}$

$$\begin{aligned} \sigma &= \frac{R^2}{4} \int d\Omega = \frac{R^2}{4} \times 4\pi = \pi R^2 \\ &= \text{cross-section area of sphere} \end{aligned}$$

Now, that is all completely classical

But ideas carry over into quantum mechanics

Quantum Scattering

Look for solutions to Schr. Eqn with form

$$\psi(r) \rightarrow A \left\{ e^{ikr} + f(\theta) \frac{e^{ikr}}{r} \right\} \quad (\text{at large } r)$$

↑ incident plane wave
↑ scattered spherical wave

Here $k = \frac{\sqrt{2mE}}{\hbar}$ as usual

Relate this to $D(\theta)$: μsr $D(\theta) = \frac{1}{L} \frac{dN}{d\Omega}$

$$dN = \frac{\# \text{ of particles}}{\text{time}} \text{ scattered into } d\Omega$$

$$= |A_{\text{scat}}|^2 \times \overset{\substack{\uparrow \\ \text{area}}}{r^2 d\Omega} \times \overset{\substack{\uparrow \\ \text{particle velocity}}}{v}$$

↑ particles per volume

$$= |A|^2 |f(\theta)|^2 \frac{1}{r^2} r^2 d\Omega v$$

$L = \text{incident particles/area} \cdot \text{time}$

$$= |A_{\text{inc}}|^2 \cdot v = |A|^2 v$$

$$\text{So } \frac{dN}{L d\Omega} = \frac{|A|^2 |f(\theta)|^2 d\Omega v}{|A|^2 v d\Omega} = |f(\theta)|^2$$

$$\text{So } \boxed{D(\theta) = |f(\theta)|^2}$$

Get $f(\theta)$ by solving Schr. Eqn.