

## Lecture 3 - Atoms & Solids

Last time, introduced spin-statistics theorem:

$$\begin{aligned} \text{integer spin} &= \text{boson} = +\text{exchange} \\ \text{half-integer spin} &= \text{fermion} = -\text{exchange} \end{aligned}$$

Consequences:

See  
supplement  
online

- Two fermions can't occupy same single particle state (including spin)
- Two bosons can.
- Two fermions in different states more likely to be found far apart
- Two bosons more likely to be found close together
- Symmetry doesn't matter if single particle states have no overlap.

Sometimes think about this using "exchange forces"

"Quantum" force that attracts bosons & repels fermions

Force doesn't come from Hamiltonian, but it is real  
 $\rightarrow$  takes a lot of energy to push two fermions together

Spin effects:

With spin, typically have

$$2t(\vec{r}_1, \vec{s}_1; \vec{r}_2, \vec{s}_2) = 2t(r_1, r_2) X(s_1, s_2) \quad \text{- spin decoupled}$$

$$\rho 2t(\vec{r}_1, \vec{s}_1; \vec{r}_2, \vec{s}_2) = 2t(\vec{r}_2, \vec{s}_2; \vec{r}_1, \vec{s}_1)$$

Need overall  $\psi$  to have correct exchange

singlet:

$$\chi = \frac{1}{\sqrt{2}}(\uparrow\downarrow - \downarrow\uparrow)$$

Total  $S=0$

For instance, if  $\chi(s_1, s_2) =$  spin singlet  
then  $\chi(s_2, s_1) = -\chi(s_1, s_2)$

Then for fermions, need  $\psi(r_1, r_2)$  symmetric  
bosons need antisymmetric

Last point:

How do we symmetrize wave function for more  
than two particles?

Say  $N$  states,  $2t_a, 2t_b, \dots 2t_n$

$N$  particles  $r_1, r_2, \dots r_N$

For bosons, just sum all permutations of states  
of particles:

$$\psi(r_1, r_2, \dots, r_N) = A [ 2t_a(r_1) 2t_b(r_2) \dots 2t_n(r_N) \\ + 2t_b(r_1) 2t_a(r_2) \dots 2t_n(r_N) \\ + \dots \text{all other combinations} ]$$

For fermions harder, need to keep track of  
signs. Can use "Slater determinant":

$$\psi(r_1, \dots, r_N) = \begin{vmatrix} 2t_a(r_1) & 2t_b(r_1) & \dots & 2t_n(r_1) \\ 2t_a(r_2) & 2t_b(r_2) & \dots & 2t_n(r_2) \\ \vdots & & & \\ 2t_a(r_N) & 2t_b(r_N) & \dots & 2t_n(r_N) \end{vmatrix}$$

Evaluate with normal rules for determinants

Can use Pauli exclusion principle to understand multi-electron atoms.

For instance, He has 2 electrons.

To first approximation, say they don't interact.

Then each electron in a single particle hydrogen state

Lowest energy: both in  $1s$  state  
only possible for spin singlet

First excited state, one  $e^-$  in  $2s$  state

Could now have spin singlet or triplet,  
which has lower energy?

If spin singlet, spatial  $\Psi$  is symmetric  
 $\Rightarrow$  electrons like to be close

Spin triplet, spatial part is antisymmetric  
 $\Rightarrow$  electrons further apart.

But since electrons both negative charge,  
cost EM energy to have close together

$\Rightarrow$  spin triplet state is lower. - correct

Section 5.2 discusses in more detail,  
explains spectroscopic notation

I'll let you read it,

Now talk about solids, particularly metals

Solid = bunch of atoms mashed together

Metal = each atom gives up a few electrons free to drift through material

That's why metals conduct

Put some electrons in on one side of wire

Take some out on other

Complicated system:

Electrons interact with ionic cores & each other

Simplest model: say electrons are free particles in a box

Turns out to work tolerably well!

Develop model: Free Electron Gas

$$\text{Box potential } V(x, y, z) = \begin{cases} 0 & 0 < x < L_x \\ 0 & 0 < y < L_y \\ 0 & 0 < z < L_z \\ \infty & \text{else} \end{cases}$$

(can write  $V = V_x(x) + V_y(y) + V_z(z)$   
each 1D infinite square well)

So wave function separates

$$\psi(x, y, z) = X(x) Y(y) Z(z)$$

We know square well solutions:

$$X(x) = \sqrt{\frac{2}{\ell_x}} \sin \frac{n_x \pi}{\ell_x} x$$

integer  $n_x > 0$

So 3D wavefunctions are

$$\psi_{n_x n_y n_z}(x, y, z) = \sqrt{\frac{8}{V}} \sin\left(\frac{n_x \pi}{\ell_x} x\right) \sin\left(\frac{n_y \pi}{\ell_y} y\right) \sin\left(\frac{n_z \pi}{\ell_z} z\right)$$

$V = \text{volume } \ell_x \ell_y \ell_z$

with energies

$$E_{n_x n_y n_z} = \frac{\hbar^2 \pi^2}{2m} \left( \frac{n_x^2}{\ell_x^2} + \frac{n_y^2}{\ell_y^2} + \frac{n_z^2}{\ell_z^2} \right)$$

$$= \frac{\hbar^2 k^2}{2m}$$

for  $\vec{k} = \left( \frac{\pi n_x}{\ell_x}, \frac{\pi n_y}{\ell_y}, \frac{\pi n_z}{\ell_z} \right)$

But that's not the whole story...

even free electrons satisfy exchange requirement

Only two electrons in each state  $(n_x, n_y, n_z)$   
 [those two  $e^-$ 's are in spin singlet]

So if we have a lot of  $e^-$ 's, we'll occupy  
 a lot of states.

→ What is total energy of electron gcs?

Use geometrical construction

Allowed wave vectors  $\vec{k}$  form grid in "k-space"

