

Gravitation and Cosmology

Lecture 18: Grad, Div, Curl, and all that...

Grad, Div, Curl, and all that^{1/4}

Reading: Mathews & Walker, *Mathematical Methods of Physics*, ch. 15.
 S. Weinberg, *Gravitation and Cosmology*, ch. 3 & 4.
 Ohanian and Ruffini, *Gravitation and Spacetime*, ch. 6 & 7.
 McConnell, *Applications of Tensor Analysis*, Ch. 12.

The generally covariant differential operators

The gradient operator is obvious, and we have already derived it. If $\varphi(x)$ is a scalar field, then

$$\varphi_{, \mu} \stackrel{df}{=} \frac{\partial \varphi}{\partial x^\mu} \quad (18.1)$$

in fact transforms as a covariant vector under general coordinate transformations.

Next, consider the covariant curl defined by

$$\begin{aligned} \text{curl}_V A_\mu &\stackrel{df}{=} A_{\mu; \nu} - A_{\nu; \mu} = A_{\mu, \nu} - A_{\nu, \mu} - \left[\left\{ \begin{matrix} \lambda \\ \mu\nu \end{matrix} \right\} - \left\{ \begin{matrix} \lambda \\ \nu\mu \end{matrix} \right\} \right] A_\lambda \\ &\equiv A_{\mu, \nu} - A_{\nu, \mu} . \end{aligned} \quad (18.2)$$

Finally, the divergence should generalize the flat-space result $\text{div } V \stackrel{df}{=} V^\mu_{, \mu}$:

$$\text{div } V \stackrel{df}{=} V^\mu_{; \mu} \equiv V^\mu_{, \mu} + \left\{ \begin{matrix} \mu \\ \mu\sigma \end{matrix} \right\} V^\sigma . \quad (18.2)$$

However,

$$\left\{ \begin{matrix} \mu \\ \mu\sigma \end{matrix} \right\} = \frac{1}{2} g^{\mu\lambda} [g_{\mu\lambda, \sigma} + g_{\lambda\sigma, \mu} - g_{\mu\sigma, \lambda}] \equiv \frac{1}{2} g^{\mu\lambda} g_{\mu\lambda, \sigma} \quad (18.3)$$

where we have used the antisymmetry in $\mu\lambda$ of the terms $g_{\lambda\sigma, \mu} - g_{\mu\sigma, \lambda}$ to drop them after contraction with $g^{\mu\lambda}$.

Now consider $g_{\mu\nu}$ as a matrix G , with $dG = g_{\mu\nu, \sigma} dx^\sigma$; then

$$\left\{ \begin{matrix} \mu \\ \mu\sigma \end{matrix} \right\} dx^\sigma = \frac{1}{2} \text{Trace} [G^{-1} dG] . \quad (18.4)$$

Let us call $\det(G) = -g$ (this is a standard notation, with the $-$ sign introduced to make g positive); then

$$-g = e^{\text{Trace} [\log(G)]} \quad (18.5)$$

Gravitation and Cosmology

The divergence theorem

(Eq. 18.5 is far from obvious, useful, and *worth remembering*. The proof is given below. For now, just believe it!)

Thus

$$\log(-g) = \text{Tr} [\log(G)]$$

and

$$\text{Tr} [\log(G + dG)] = \text{Tr} [\log(G)] + \text{Tr} [\log(1 + G^{-1}dG)],$$

hence

$$\log(-g - dg) = \log(-g) + \text{Tr} [G^{-1}dG]$$

so that

$$\left\{ \begin{matrix} \mu \\ \mu\sigma \end{matrix} \right\} dx^\sigma = \frac{1}{2} d[\log(-g)] \equiv d[\log(\sqrt{-g})]$$

or (we can now drop the $-$ sign from $\sqrt{-g}$)

$$\left\{ \begin{matrix} \mu \\ \mu\sigma \end{matrix} \right\} = \partial_\sigma \log(\sqrt{g}) = \frac{1}{\sqrt{g}} \partial_\sigma \sqrt{g}. \quad (18.6)$$

“So what?” you may say, “I’ve got my own troubles.” Here’s what:

Remember Eq. 18.2? Now we may write

$$V^\mu{}_{;\mu} = V^\mu{}_{,\mu} + \frac{1}{\sqrt{g}} V^\mu \partial_\mu \sqrt{g} \equiv \frac{1}{\sqrt{g}} \partial_\mu (V^\mu \sqrt{g}). \quad (18.7)$$

That is, the expression for the covariant divergence is charmingly simple.

The divergence theorem

Under coordinate transformations, the volume element changes like

$$d^n x \rightarrow d^n \bar{x} \det \left(\frac{\partial x}{\partial \bar{x}} \right) \quad (18.8)$$

where $\det \left(\frac{\partial x}{\partial \bar{x}} \right)$ is the Jacobian of the transformation.

But

$$\bar{g}_{\mu\nu} = g_{\kappa\lambda} \frac{\partial x^\kappa}{\partial \bar{x}^\mu} \frac{\partial x^\lambda}{\partial \bar{x}^\nu} \quad (18.9)$$

so, using a well-known property of determinants of matrix-products,

$$\det(AB) = \det(A) \det(B), \quad (18.10)$$

we find

$$\bar{g} = g \left[\det \left(\frac{\partial x}{\partial \bar{x}} \right) \right]^2 \quad (18.11)$$

i.e.

$$\sqrt{\bar{g}} d^n x = \sqrt{g} d^n \bar{x}. \quad (18.12)$$

Gravitation and Cosmology

Lecture 18: Grad, Div, Curl, and all that...

In other words, $\sqrt{g} d^n x$ is the invariant volume element in a general space of n dimensions.

Now we can transform the volume integral of a divergence into an integral of a vector over a (normal) hypersurface.

$$\int \operatorname{div} V \sqrt{g} d^n x = \int V^\mu{}_{;\mu} \sqrt{g} d^n x = \int d^n x \partial_\mu (\sqrt{g} V^\mu) \equiv \int dS_\mu V^\mu \sqrt{g}. \quad (18.13)$$

Proof of the theorem about determinants:

We want to prove that for some matrix G ,

$$\log [\det(G)] = \operatorname{Tr} [\log(G)].$$

Now this is obvious if the matrix is diagonalizable, with eigenvalues g_k since then

$$\log [\det(G)] = \log \left[\prod_{k=1}^N g_k \right] \equiv \sum_{k=1}^N \log(g_k)$$

and

$$\operatorname{Tr} [\log(G)] = \sum_{k=1}^N \log(g_k).$$

We are concerned to prove the theorem more generally. First, it had better be true that the matrix

$$A = \overset{df}{\log}(G)$$

exists (that is, it can be defined, the matrix has no zero eigenvalues, *etc. etc.*). Assuming this is the case, let

$$G(\lambda) = \overset{df}{e^{\lambda A}}, \quad G(1) = e^A = G.$$

Now let us define

$$\begin{aligned} d \log [\det(G(\lambda))] &= \overset{df}{\log} [\det(G(\lambda + d\lambda))] - \overset{df}{\log} [\det(G(\lambda))] \\ &= \log [\det(G + dG)] - \log [\det(G)] \\ &= \log [\det(G(1 + G^{-1}dG))] - \log [\det(G)] \\ &= \log [\det(1 + G^{-1}dG)]. \end{aligned}$$

Now let us compute the last term:

Gravitation and Cosmology

Proof of the theorem about determinants:

$$\begin{aligned} \det(1 + G^{-1}dG) &= \begin{vmatrix} 1 + (G^{-1}dG)_{11} & (G^{-1}dG)_{12} & (G^{-1}dG)_{13} & \dots \\ (G^{-1}dG)_{21} & 1 + (G^{-1}dG)_{22} & (G^{-1}dG)_{23} & \dots \\ (G^{-1}dG)_{31} & (G^{-1}dG)_{32} & 1 + (G^{-1}dG)_{33} & \dots \\ \dots & \dots & \dots & \dots \end{vmatrix} \\ &= \prod_k \left[1 + (G^{-1}dG)_{kk} \right] + O((G^{-1}dG)^2) \\ &\approx 1 + \sum_k (G^{-1}dG)_{kk} \equiv 1 + \text{Tr}(G^{-1}dG). \end{aligned}$$

To this same order, then,

$$d \log [\det(G(\lambda))] = \log \left[1 + \text{Tr}(G^{-1}dG) \right] = \text{Tr}(G^{-1}dG)$$

However, since $G(\lambda) = e^{\lambda A}$, clearly

$$dG(\lambda) = A e^{\lambda A} d\lambda$$

and thus

$$d \log [\det(G(\lambda))] = \text{Tr}(G^{-1}dG) = \text{Tr}(e^{-\lambda A} A e^{\lambda A}) d\lambda \equiv \text{Tr}(A) d\lambda,$$

giving, by direct integration,

$$\log [\det(G(\lambda))] = \lambda \text{Tr}(A) + \text{constant}.$$

Since both sides must vanish when $\lambda = 0$, the constant is zero, giving at last (with $\lambda = 1$)

$$\det(G) = \exp [\text{Tr}(A)] = \exp [\text{Tr}(\log(G))].$$