

1. Find the Fourier transform of the function $f(x)$ where:

$$f(x) = \begin{cases} 0 & (\text{if } x < -a/2) \\ e^{i\beta x} & (\text{if } -a/2 < x < a/2) \\ 0 & (\text{if } a/2 < x < b - a/2) \\ e^{i\beta x} & (\text{if } b - a/2 < x < b + a/2) \\ 0 & (\text{if } b + a/2 < x) \end{cases}$$

Here a and b are positive real constants with $b > a$.

2. Calculate the convolution

$$g(t) = \int_{-\infty}^{\infty} f_1(T)f_2(t-T)dT$$

for

$$f_1(t) = \begin{cases} 1 & (\text{if } -\tau_1 < t < \tau_1) \\ 0 & (\text{otherwise}) \end{cases}$$

$$f_2(t) = \begin{cases} 1 & (\text{if } -\tau_2 < t < \tau_2) \\ 0 & (\text{otherwise}) \end{cases}$$

where $\tau_1 > \tau_2$. Sketch a plot of $g(t)$, and note its full-width at half-maximum.

Hint: you will need to separately consider the cases (a) $t < -(\tau_1 + \tau_2)$; (b) $-(\tau_1 + \tau_2) < t < \tau_2 - \tau_1$; (c) $\tau_2 - \tau_1 < t < \tau_1 - \tau_2$; (d) $\tau_1 - \tau_2 < t < \tau_1 + \tau_2$; (e) $\tau_1 + \tau_2 < t$. For each of these cases, it helps to draw a picture of $f_1(T)$ and $f_2(t-T)$ as functions of T .

3. For each of the following aperture functions, calculate the propagating wave $E(x, y, z)$, assuming light with wave number k :

(a) $A(x, y) = e^{iky/4}$

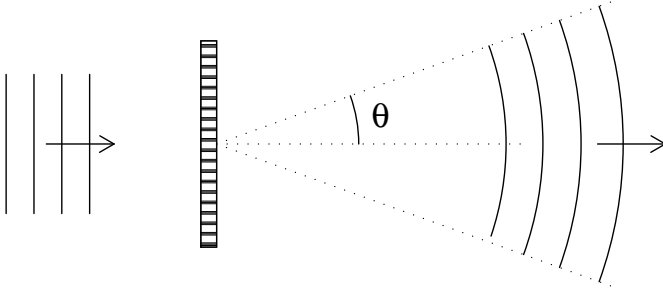
(b) $A(x, y) = e^{-ik(x+y)/2}$

(c) $A(x, y) = e^{ik(x+y)}$

(d) $A(x, y) = \sin^2(ky/4)$

4. Suppose that light of wavelength $\lambda = 633$ nm can be described in the $z = 0$ plane by an aperture function $A(x, y)$. It is known that $A(x, y)$ has a Fourier spectrum which is zero for $\sqrt{k_x^2 + k_y^2} > 10^6 \text{ m}^{-1}$. Show that the Fraunhofer diffraction pattern will be contained entirely within a cone and calculate the cone angle θ . (See picture on reverse.)

Problem 4:



5. The Fourier method we have developed for diffraction can be applied to other problems as well. For instance, consider a pulse of light $E(z, t)$ incident on a dispersive medium with index $n(\omega)$ and length d . The front face of the medium is at $z = 0$, and the incident pulse is

$$E(z = 0, t) \equiv A(t) = E_0 e^{-t^2/\tau^2} e^{-i\omega_0 t}$$

for pulse duration τ and carrier frequency ω_0 . We can ask: what is the form of the pulse when it exits the medium?

Using the Fourier transform, the incident pulse can be written as a sum of harmonic functions $e^{-i\omega t}$. Each of these components propagates through the medium as $e^{i(kz - \omega t)}$ and thus exits after acquiring a phase shift $\mathcal{H}(\omega) = e^{ikd}$ for $k = k(\omega) = n\omega/c$. Symbolically,

$$(\text{input}) \ e^{-i\omega t} \rightarrow e^{-i\omega t} e^{ikd} \ (\text{output})$$

By resumming the transmitted harmonic functions, the complete transmitted pulse can be calculated.

Suppose the wavenumber k can be approximated as

$$k(\omega_0 + \Delta) = k_0 + k_1 \Delta + \frac{1}{2} k_2 \Delta^2$$

where $\Delta = \omega - \omega_0$, $k_0 = k(\omega_0)$, $k_1 = dk/d\omega$, and $k_2 = d^2k/d\omega^2$. (Take k_0 , k_1 and k_2 as givens, don't try to somehow calculate them.) Use the procedure outlined above to calculate the electric field of the transmitted pulse. In particular, find:

- (a) the time at which the peak of the pulse exits the medium, and
- (b) the duration of the pulse when it exits the medium (ie, the new value of τ).

Hint: The Gaussian beam example from lecture 15, slides 32–35 illustrates a similar calculation.