

1. Rate equations are

$$\frac{dN_1}{dt} = -RN_1 + \Gamma_2 N_2 + RN_3 \quad (1)$$

$$\frac{dN_2}{dt} = +\Gamma_3 N_3 - \Gamma_2 N_2 \quad (2)$$

$$\frac{dN_3}{dt} = +RN_1 - \Gamma_3 N_3 - RN_3 \quad (3)$$

Solve in steady state $\frac{d}{dt} \rightarrow 0$

$$\text{from (2), } N_3 = \frac{\Gamma_2}{\Gamma_3} N_2$$

$$\text{from (3), } RN_1 - (\Gamma_3 + R) \frac{\Gamma_2}{\Gamma_3} N_2 = 0$$

$$\text{So } N_2 = \frac{\Gamma_3}{\Gamma_2} \frac{R}{\Gamma_3 + R} N_1$$

Get inversion if $N_2 > N_1$, so need

$$\frac{\Gamma_3}{\Gamma_2} \frac{R}{\Gamma_3 + R} > 1$$

$$\Gamma_3 R > \Gamma_2 \Gamma_3 + \Gamma_2 R$$

$$(\Gamma_3 - \Gamma_2) R > \Gamma_2 \Gamma_3$$

or

$$R > \frac{\Gamma_2 \Gamma_3}{\Gamma_3 - \Gamma_2}$$

(2)

2. Gain coefficient is

$$\gamma = -\frac{1}{8\pi} \lambda_{12}^2 \Delta N \frac{g(v)}{t_s}$$

$$\text{here } g(v) \approx \frac{1}{\Delta v_{12}} = \frac{1}{10 \text{ GHz}}$$

Need ΔN : from rate equations in class,

$$-\Delta N = RN \left(\frac{1}{n_2} - \frac{g_2}{g_1} \frac{1}{n_1} \right)$$

where R is pump rate on $0 \rightarrow 3$ transition,
and we assume $R \ll \Gamma$'s

For optical pumping, have

$$R = \frac{1}{8\pi} \lambda_{03}^2 \frac{I}{hv_{03}} \frac{g_3}{g_0} \frac{g(v)}{t_s}$$

$$\text{with } g(v) \approx \frac{1}{\Delta v_{03}} = \frac{1}{100 \text{ GHz}}$$

$= 1140 \text{ s}^{-1}$ which is indeed small
compared to Γ 's

$$\begin{aligned} \text{So, } -\Delta N &= (1140 \text{ s}^{-1}) \left(30 \text{ ns} - \frac{1}{2} 5 \text{ ns} \right) N \\ &= 3.1 \times 10^{-5} N \\ &= 3.1 \times 10^{18} \text{ m}^{-3} \end{aligned}$$

and

$$\begin{aligned} \gamma &= \frac{1}{8\pi} \left(\frac{800 \text{ nm}}{1.5} \right)^2 \left(3.1 \times 10^{18} \text{ m}^{-3} \right) \frac{1}{10 \text{ GHz} \cdot 30 \text{ ns}} \\ &= 117 \text{ m}^{-1} = \boxed{1.2 \cdot \text{cm}^{-1}} \end{aligned}$$

(4)

3.

Total loss per pass is $1\% + \frac{1}{2} \times 4\% = 3\%$

available gain is $0.2 \text{ m}^{-1} \times 20 \text{ cm} = 0.04$,

so laser will oscillate

In steady state,

$$\gamma = \frac{\gamma_0}{1 + \frac{I}{I_s}} = 0.03$$

$$\text{So } 1 + \frac{I}{I_s} = \frac{4}{3}$$

$$I = \frac{1}{3} I_s$$

have $I_{\text{set}} = \frac{4\pi h\nu}{\lambda^2 g(\nu)}$ at peak, $g(\nu) = 2T_{12} = 4t_s$

$$I_{\text{set}} = \frac{4\pi h\nu}{\lambda^2 t_s}$$

$$= 229 \frac{\text{W}}{\text{m}^2}$$

$$\text{So } I_{\text{cav}} = 76 \frac{\text{W}}{\text{m}^2} = I \text{ in cavity}$$

$$\text{and } I_{\text{out}} = T I_{\text{cav}} = 0.04 \times I_{\text{cav}} \\ = 3.5 \frac{\text{W}}{\text{m}^2}$$

To get power, need beam area.

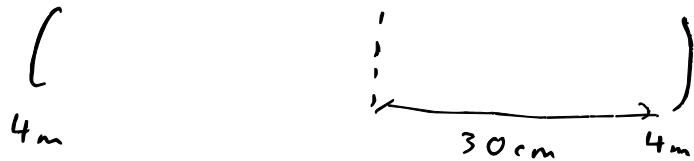
Analyze cavity:



3. (cont)

(5)

This looks like half of symmetric cavity



$$\text{From Yariv 7.1-9, } w_{\text{mirror}} = \sqrt{\frac{\lambda \ell}{2\pi}} \left[\frac{2R^2}{\ell(R-\frac{\ell}{2})} \right]^{1/4}$$

for $\ell = 60 \text{ cm}$

and

$$w_0 = \sqrt{\frac{2}{\pi}} \left[\frac{\ell}{2} \left(R - \frac{\ell}{2} \right) \right]^{1/4}$$

$$w_{\text{mirror}} = 370 \mu\text{m}$$

$$w_0 = 342 \mu\text{m}$$

So, beam waist is reasonably uniform $w \approx 350 \mu\text{m}$

$$\text{Then } P \propto \pi w^2 I$$

$$P = 1.3 \mu\text{W}$$

$$\begin{aligned} \text{For optimum coupling, } T &= -L_i + \sqrt{s_0 L_i} && (\text{L}_i \text{ and } s_0 \text{ per round trip}) \\ &= -0.02 + \sqrt{0.02 \times 0.08} = 0.02 \end{aligned}$$

$$\text{Then } P_{\text{out}} = I_s A / (\sqrt{s_0} - \sqrt{L})^2 \quad (\text{Yariv 9.3-18})$$

$$= 230 \frac{\text{W}}{\text{m}^2} \pi w^2 / (\sqrt{0.08} - \sqrt{0.02})^2 = 1.8 \mu\text{W}$$

(6)

4.

$$\text{Have } \nu_m' = \frac{\nu_m}{1 + \chi'/2}$$

where ν_m' = mode frequency with gas present

ν_m = mode frequency w/o gas

$$\text{Here } \chi' = \frac{2(v_0 - v)}{\Delta v} \chi''$$

what is χ'' ?

$$\text{Know } \alpha_0 = k \chi''(v_0)$$

$$\text{and } \chi''(v) = \chi''(v_0) \frac{g(v)}{g(v_0)}$$

$$\text{with } g(v) = \frac{2T_2}{1 + 4\pi^2 T_2^2 (v - v_0)^2}$$

$$\Delta v = \frac{1}{\pi T_2}$$

$$\text{thus } g(v) = \frac{2}{\pi \Delta v} \frac{1}{1 + 4 \frac{(v - v_0)^2}{\Delta v^2}}$$

$$\text{and } g(v_0) = \frac{2}{\pi \Delta v}$$

Putting together,

$$\chi' = \frac{2(v_0 - v)}{\Delta v} \frac{\alpha_0}{K} \frac{1}{1 + 4 \frac{(v - v_0)^2}{\Delta v^2}}$$

4. (cont)

(7)

We're interested in v_m' when $v_m = v_0 \pm \Delta v$

Take + case:

$$v'_+ = \frac{v_0 + \frac{c}{2\ell}}{1 + \frac{1}{2}x'(v'_+)} = v_0 + \delta$$

$$= \frac{v_0 + \frac{c}{2\ell}}{1 + \frac{(-\delta)}{\Delta v} \frac{\alpha_0}{k} \left(\frac{1}{1 + 4 \frac{\delta^2}{\Delta v^2}} \right)}$$

But we know $\Delta v \gg \frac{c}{2\ell}$, so expect $\Delta v \gg \delta$

$$\text{So, } v'_+ = \frac{v_0 + \frac{c}{2\ell}}{1 - \frac{\delta}{\Delta v} \frac{\alpha_0}{k}} = v_0 + \delta$$

$$v_0 + \frac{c}{2\ell} = (v_0 + \delta) \left(1 - \frac{\delta}{\Delta v} \frac{\alpha_0}{k} \right)$$

$$\frac{c}{2\ell} = \delta - \frac{\delta v_0}{\Delta v} \frac{\alpha_0}{k}$$

$$= \delta \left(1 - \frac{v_0 \alpha_0}{\Delta v k} \right)$$

$$\text{So } \delta = \frac{c}{2\ell} \left(1 + \frac{v_0 \alpha_0}{\Delta v k} \right)$$

$$\text{note } k = \frac{\omega}{c} = \frac{2\pi v_0}{c}$$

4. (cont)

(8)

$$\delta = \Delta v_L \left(1 + \frac{\alpha_0 C}{2\pi \Delta v} \right)$$

Check ... is $\frac{\alpha_0 C}{2\pi \Delta v} \ll 1$?

$$= \frac{\Delta v_L}{\Delta v} \frac{\alpha_0 C}{2\pi \Delta v_L}$$

$$= \frac{\Delta v_L}{\Delta v} \frac{\alpha_0 C}{2\pi \frac{C}{2L}}$$

$$= \frac{\Delta v_L}{\Delta v} \underbrace{\frac{1}{\pi}}_{\ll 1} \underbrace{\alpha_0 L}_{\approx 1}$$

So, shift is small, as we assumed.

Note shift is away from v_0 ... looks like

"frequency pushing" here

(difference is gain vs. absorption)