

1. a) From HW 1, have round trip matrix

$$\begin{bmatrix} -1.256 & 1.4508 \text{ m} \\ -1.8667 \text{ m}^{-1} & 1.36 \end{bmatrix}$$

for round trip starting just before mirror M_1 .

$$\begin{aligned} \text{The at this point, } q &= \frac{1}{C} \left[\frac{A-D}{2} \pm i \sqrt{1 - \left(\frac{A+D}{2}\right)^2} \right] \\ &= 0.7 \text{ m} + i 0.535 \text{ m} \end{aligned}$$

So for leg 3-4-1, have focus 700 mm behind M_1 , or 130 mm behind M_4

[Note mirror M_4 has no effect, since $R_4 = \infty$]

$$\text{For this focus, } w_0 = \sqrt{\frac{\lambda z_0}{\pi}} = \text{span style="border: 1px solid black; padding: 2px;">370 \mu\text{m}}$$

For leg 1-2, just propagate q through mirror M_1 :

$$\text{matrix} = \begin{bmatrix} 1 & 0 \\ \frac{2}{R_1} & 1 \end{bmatrix} \quad R_1 = -0.1 \text{ m}$$

$$\text{So } q' = \frac{q}{\frac{2}{R_1} q + 1} = -52.3 \text{ mm} + i 1.887 \text{ mm}$$

So, get focus 52.3 mm in front of M_1 .

beam waist $w_0 = \sqrt{\frac{\lambda z_0}{\pi}} = 21.9 \mu\text{m}$

Finally, check leg 2-3:

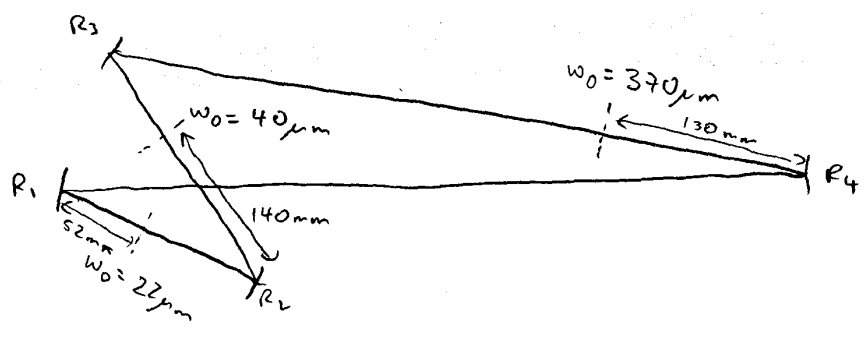
Use matrix $M = \begin{bmatrix} 1 & 0 \\ -\frac{2}{0.1} & 1 \end{bmatrix} \begin{bmatrix} 1 & 0.13 \\ 0 & 1 \end{bmatrix}$
 $= \begin{bmatrix} 1 & 0.13 \text{ m} \\ -20 \text{ m}^{-1} & -1.6 \end{bmatrix}$

$q' = \frac{q + 0.13}{-20q - 1.6}$ for $q = -5.23 \times 10^{-3} + i 1.887 \times 10^{-3} \text{ m}$

$= -139.8 + i 6.12 \text{ mm}$

Focus is 139.8 mm past M_2 , with waist $= \sqrt{\frac{\lambda z_0}{\pi}} = 39.5 \mu\text{m}$

Picture:



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b) Output beam is diverging from focus with

$$w_0 = 370 \mu\text{m}$$

$$\text{So, has divergence angle } \Theta = \frac{\lambda}{\pi w_0} = \boxed{0.7 \text{ mrad}}$$

$$\text{Free spectral range } \nu_F = \frac{c}{L}$$

$$L = 130 + 220 + 660 + 570 \text{ mm} = 1.58 \text{ m}$$

$$\boxed{\nu_F = 190 \text{ MHz}}$$

Gouy phase

$$\text{leg 1-2: } z_1 = -52.3 \text{ mm} \quad z_2 = +77.7 \text{ mm}$$

$$z_0 = 1.887 \text{ mm}$$

$$\Delta S_{12} = \tan^{-1} \frac{z_2}{z_0} - \tan^{-1} \frac{z_1}{z_0}$$

$$= 1.546 + 1.535 = 3.081 \text{ rad}$$

$$\text{leg 2-3: } z_1 = -139.8 \text{ mm} \quad z_2 = 80.2 \text{ mm}$$

$$z_0 = 6.12 \text{ mm}$$

$$\Delta S_{23} = \tan^{-1} \frac{80.2}{6.12} + \tan^{-1} \frac{139.8}{6.12}$$

$$= 3.022 \text{ rad}$$

$$\text{leg 3-1: } z_1 = -530 \text{ mm} \quad z_2 = 700 \text{ mm}$$

$$z_0 = 535 \text{ mm}$$

$$\Delta S_{31} = \tan^{-1} \frac{700}{535} + \tan^{-1} \frac{530}{535}$$

$$= 1.699 \text{ rad}$$

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$$\text{So } \Delta S_{\text{Tot}} = \Delta S_{12} + \Delta S_{23} + \Delta S_{31}$$

$$\Delta S = 7.802 \text{ rad}$$

After one pass through cavity, intensity is reduced by factor $(0.995)^3 \times 0.95 = 0.9358$

$$\text{So, loss per pass} = 0.0642 = \Gamma$$

$$\text{Then } \delta\nu = \nu_F \frac{\Gamma}{2\pi} = 1.94 \text{ MHz}$$

$$\text{Finesse } \mathcal{F} = \frac{\nu_F}{\delta\nu} = \frac{2\pi}{\Gamma} = 97.8$$

$$Q = \frac{\nu}{\delta\nu} = \frac{c}{\lambda \delta\nu} = 1.93 \times 10^8$$

$$T_p = \frac{1}{2\pi \delta\nu} = 82 \text{ ns}$$

$$\alpha = \frac{1}{L} \ln \frac{1}{1-\Gamma} = 0.042 \text{ m}^{-1}$$

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4. a)

$$m\ddot{x} + b\dot{x} + kx = 0$$

Rewrite: $\ddot{x} + \gamma\dot{x} + \omega_0^2 x = 0$

$$\gamma = \frac{b}{m}$$

$$\omega_0^2 = \frac{k}{m}$$

Plus in $x(t)$:

$$\dot{x} = A\left(-\frac{\gamma}{2}\right) e^{-\gamma t/2} \cos(\bar{\omega}t + \phi) - A\bar{\omega} e^{-\gamma t/2} \sin(\bar{\omega}t + \phi)$$

$$\ddot{x} = A\left(\frac{\gamma^2}{4}\right) e^{-\gamma t/2} \cos(\bar{\omega}t + \phi) + 2A\frac{\gamma}{2}\bar{\omega} e^{-\gamma t/2} \sin(\bar{\omega}t + \phi) - A\bar{\omega}^2 e^{-\gamma t/2} \cos(\bar{\omega}t + \phi)$$

So have equation:

$$A e^{-\gamma t/2} \left[\left(\frac{\gamma^2}{4} - \bar{\omega}^2\right) \cos(\bar{\omega}t + \phi) + \gamma\bar{\omega} \sin(\bar{\omega}t + \phi) \right] =$$

$$+ \gamma A e^{-\gamma t/2} \left[-\frac{\gamma}{2} \cos(\bar{\omega}t + \phi) - \bar{\omega} \sin(\bar{\omega}t + \phi) \right]$$

$$+ \omega_0^2 A e^{-\gamma t/2} \cos(\bar{\omega}t + \phi) = 0$$

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So $Ae^{-\gamma t/2}$ cancels out, set

$$\left(\frac{\gamma^2}{4} - \bar{\omega}^2 + \omega_0^2 - \frac{\gamma}{2}\sigma\right) \cos(\bar{\omega}t + \phi) \\ + (\gamma\bar{\omega} - \sigma\bar{\omega}) \sin(\bar{\omega}t + \phi) = 0$$

Coefficients of sine and cosine must vanish, so

$$\boxed{\gamma = \sigma = b/m}$$

and

$$\bar{\omega}^2 = \omega_0^2 + \frac{\gamma^2}{4} - \frac{\gamma^2}{2} = \omega_0^2 - \frac{\gamma^2}{4}$$

$$\boxed{\bar{\omega} = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}}$$

b) Then

$$E = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \omega_0^2 x^2$$

$$= \frac{1}{2} m A^2 e^{-\gamma t} \left[-\frac{\gamma}{2} \cos(\bar{\omega}t + \phi) - \bar{\omega} \sin(\bar{\omega}t + \phi) \right]^2 \\ + \frac{1}{2} m A^2 e^{-\gamma t} \omega_0^2 \cos^2(\bar{\omega}t + \phi)$$

$$= \frac{1}{2} m A^2 e^{-\gamma t} \left[\frac{\gamma^2}{4} \cos^2(\bar{\omega}t + \phi) + \gamma\bar{\omega} \cos(\bar{\omega}t + \phi) \sin(\bar{\omega}t + \phi) \right. \\ \left. + \bar{\omega}^2 \sin^2(\bar{\omega}t + \phi) + \omega_0^2 \cos^2(\bar{\omega}t + \phi) \right]$$

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Use $\bar{\omega}^2 = \omega_0^2 - \frac{\gamma^2}{4}$, and trig identities

$$E = \frac{1}{2} m A^2 e^{-\gamma t} \left[\omega_0^2 + \frac{\gamma^2}{4} \cos 2(\bar{\omega}t + \phi) + \frac{\gamma}{2} \bar{\omega} \sin 2(\bar{\omega}t + \phi) \right]$$

c) If $\gamma \ll \omega$, then $e^{-\gamma t}$ doesn't change much during one cycle, so \cos & \sin term average out

$$\langle E \rangle \approx \frac{1}{2} m \omega_0^2 A^2 e^{-\gamma t}$$

$$\frac{d\langle E \rangle}{dt} = -\frac{\gamma}{2} m \omega_0^2 A^2 e^{-\gamma t}$$

So $Q' = \frac{\bar{\omega}}{\gamma}$

d) If $F = F_0 e^{i\omega t}$, have

$$\ddot{x} + \gamma \dot{x} + \omega_0^2 x = f_0 e^{i\omega t} \quad f_0 = \frac{F_0}{m}$$

Try solution $x = B e^{i\omega t}$

Then

$$B(-\omega^2) + \gamma B(i\omega) + \omega_0^2 B = f_0$$

$$B = \frac{f_0}{\omega_0^2 - \omega^2 + i\gamma\omega}$$

$$e) |x|^2 = |\dot{x}|^2 = \frac{f_0^2}{(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2}$$

If $\gamma \ll \omega_0$, then resonance occurs very close to $\omega = \omega_0$

$$\text{So simplify } \omega_0^2 - \omega^2 = (\omega_0 + \omega)(\omega_0 - \omega) \\ \approx 2\omega_0(\omega_0 - \omega)$$

$$\text{and } \gamma^2 \omega^2 = \gamma^2 \omega_0^2$$

$$\text{So } |x|^2 = \frac{f_0^2}{\omega_0^2} \frac{1}{4(\omega - \omega_0)^2 + \gamma^2}$$

$$\text{Peak} = \frac{f_0^2}{\omega_0^2 \gamma^2}$$

$$\text{So half max at } 4(\omega - \omega_0)^2 + \gamma^2 = 2\gamma^2$$

$$(\omega - \omega_0)^2 = \frac{\gamma^2}{4}$$

$$\omega = \omega_0 \pm \frac{\gamma}{2}$$

$$\text{So FWHM is } \boxed{\Delta\omega = \gamma}$$

$$f) Q = \frac{\bar{\omega}}{\Delta\omega} = \frac{\bar{\omega}}{\gamma} = Q' \text{ as desired.}$$

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5. N photons have energy $\Sigma = N h \nu$

Modes of symmetric cavity have focus at center,

and
$$z_0 = \frac{1}{2} \sqrt{-d(2R+d)}$$

Total field in cavity is sum of forward & backward beams:

$$\begin{aligned} u(\rho, z) &= A_0 \frac{w_0}{w(z)} e^{-ikz} e^{-ik \frac{\rho^2}{2R(z)}} e^{-\rho^2/w(z)^2} e^{iS(z)} \\ &+ A_0 \frac{w_0}{w(z)} e^{+ikz} e^{-ik \frac{\rho^2}{2R(z)}} e^{-\rho^2/w(z)^2} e^{iS(z)} \\ &= 2 A_0 \cos kz \frac{w_0}{w(z)} e^{-ik \frac{\rho^2}{2R(z)}} e^{-\rho^2/w(z)^2} e^{iS(z)} \end{aligned}$$

So
$$I(\rho, z) = I_0 \cos^2 kz \frac{w_0^2}{w(z)^2} e^{-2\rho^2/w(z)^2}$$

$$I_0 = \text{peak intensity} = I(0,0)$$

Then total energy is

$$\begin{aligned} \mathcal{E} &= \frac{1}{c} \int_0^\infty 2\pi \rho d\rho \int_{-d/2}^{d/2} dz I(\rho, z) \\ &= \frac{I_0 w_0^2}{c} \int_{-d/2}^{d/2} dz \cos^2 kz \frac{1}{w(z)^2} \int_0^\infty 2\pi \rho e^{-2\rho^2/w(z)^2} d\rho \end{aligned}$$

From book, ρ integral is $\frac{1}{2} \pi w(z)^2$

So
$$\Sigma = \frac{I_0 w_0^2}{c} \frac{\pi}{2} \int_{-d/2}^{d/2} \cos^2 kz dz$$

$$\begin{aligned} \text{Then } \int_{-d/2}^{d/2} \cos^2 kz \, dz &= \frac{1}{2} \int_{-d/2}^{d/2} (1 + \cos 2kz) \, dz \\ &= \frac{d}{2} + \frac{1}{2k} \sin kd \end{aligned}$$

But for a resonant mode, have $kd = \pi q$
for integer q , leaving

$$\Sigma = \frac{I_0 \omega_0^2}{c} \frac{\pi d}{4} = N h \nu$$

$$I_0 = \frac{4}{\pi} \frac{N h \nu c}{\omega_0^2 d}$$

Use $\pi \omega_0^2 = z_0 \lambda$ and $\nu = \frac{c}{\lambda}$

$$I_0 = 4N \frac{h \nu^2}{z_0 d}$$

Here $\nu = \frac{c}{\lambda} = 6 \times 10^{14} \text{ s}^{-1}$

$$z_0 = \frac{1}{2} \sqrt{-d(2R+d)} = 5 \text{ cm}$$

$$I_0 = 4 \times 10^4 \frac{6.626 \times 10^{-34} \text{ J}\cdot\text{s} \cdot (6 \times 10^{14} \text{ s}^{-1})^2}{5 \text{ cm} \cdot 10 \text{ cm}}$$

$$I_0 = 0.19 \text{ W/cm}^2$$

a) Transverse mode spacing is $\frac{\Delta S}{2\pi} \nu_F$

(Book has $\frac{\Delta S}{\pi}$, but their ΔS is only half of round trip ΔS !)

From problem 2, $\Delta S = 7.802$

So $\frac{\Delta S}{2\pi} = 1.242$

and $\nu_F = 190 \text{ MHz}$

So mode spacing is 236 MHz

b) Get:

