

01/28/05

## Lecture 5

①

Wave approach to laser beam:

Approximate wave equation by paraxial Helmholtz equation:

$$\vec{E} = \hat{z} A(\vec{r}) e^{-ikz} e^{i\omega t}$$

$$\nabla_{\vec{r}}^2 A - 2ik \frac{\partial A}{\partial z} = 0$$

Found a solution

$$A(\vec{r}) = \frac{A_1}{z} e^{-ik \frac{\rho^2}{2z}}$$

here

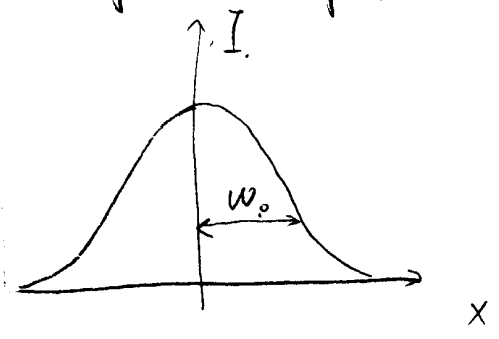
$$z = z + iz_0 \quad \text{complex}$$

$$A(\vec{r}) = \frac{A_1}{z + iz_0} e^{-\frac{kz_0 \rho^2}{z(z^2 + z_0^2)}} e^{-ik \frac{\rho^2 z}{z(z^2 + z_0^2)}}$$

$$\text{Then } I(\vec{r}) \propto |A|^2 = \frac{A_1^2}{z^2 + z_0^2} e^{-\frac{kz_0 \rho^2}{z^2 + z_0^2}}$$

$$\text{At } z = 0 = \frac{|A_1|^2}{z_0^2} e^{-\frac{k\rho^2}{z_0^2}}$$

### Gaussian function of P



Define  $1/e^2$  radius = beam waist =  $w_0$ .

Then  $I \sim e^{-2\frac{P^2}{w_0^2}}$

$$w_0^2 = z \frac{z_0}{k} = z \frac{z_0}{2\pi/\lambda} = \frac{\lambda z_0}{\pi}$$

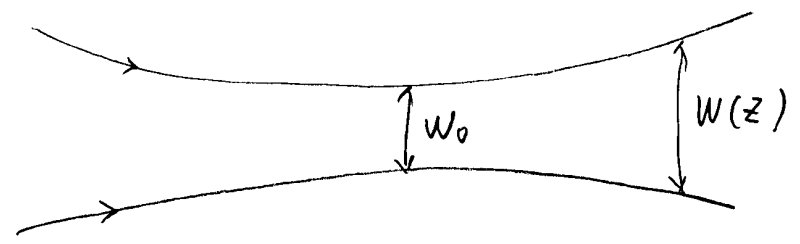
or.  $z_0 = \frac{\pi w_0^2}{\lambda}$  Rayleigh length.

At any other  $z$ , write

$$I(z) = I_0 \frac{w_0^2}{w(z)^2} e^{-\frac{2P^2}{w(z)^2}}$$

where  $w(z)^2 = w_0^2 (1 + \frac{z^2}{z_0^2})$

So intensity looks like



Describes a laser beam focussed at  $z=0$

Called a "Gaussian beam"

Good model for a laser beam.

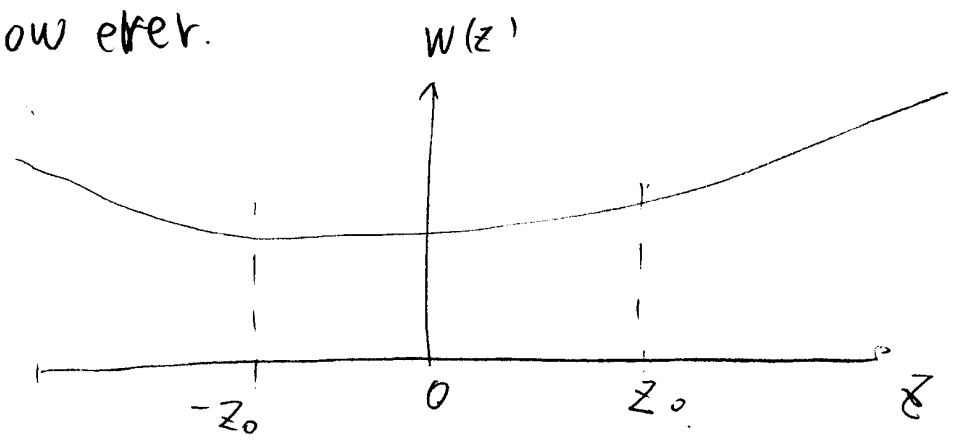
Look at properties of this beam.

a) Beam has focus at  $z=0$ .

$z < 0$ , beam is converging to focus  
 $z > 0$ , beam is diverging to focus.

Laser seems like a collimated beam w/o diverging or converging, but due to their wave nature, lasers have to have diffraction.

How ever.



looks collimated.

At  $z = -z_0$  to  $z_0$

$W(z)$  ranges from  $W_0$  to  $\sqrt{2}W_0$

Over Rayleigh range  $-z_0 < z < z_0$

beam is considered collimated.

Since  $z_0 = \frac{\pi W_0^2}{\lambda}$ , can be pretty big

$$W_0 = 1 \text{ mm} \quad \lambda = 632 \text{ nm}$$

$$\underline{z_0 = 5 \text{ m}}$$

makes it easy to not notice divergence

b) Gaussian beam defined by two parameters

(5)

Rewrite solution.

$$u(\vec{r}) = \frac{A_1}{g} e^{-ik \frac{\rho^2}{2g}} e^{-ikz}$$

Then  $\frac{1}{g} = \frac{1}{z + iz_0} = \frac{z - iz_0}{z^2 + z_0^2}$

\* Im part  $-\frac{z_0}{z^2 + z_0^2} = \frac{1}{z_0} \frac{1}{1 + \frac{z^2}{z_0^2}}$   
 $= \frac{1}{z_0^2} \frac{W_0^2}{W(z)^2}$

But  $W_0^2 = z_0^2 \pi$

so Im  $\frac{1}{g} = -\frac{\pi}{\pi W(z)^2}$

\* Also define  $\text{Re} \frac{1}{g} = \frac{z}{z^2 + z_0^2} = \frac{1}{R(z)}$

Then  $e^{-ik \frac{\rho^2}{2g}} = e^{-k \frac{\rho^2 \pi}{2\pi W(z)^2}} e^{-ik \frac{\rho^2}{2R(z)}}$   
 $= e^{-\frac{\rho^2}{W(z)^2}} e^{-ik \frac{\rho^2}{2R(z)}}$

$e^{-\frac{p^2}{w(z)^2}}$  gives amplitude, Gaussian spot. (6)

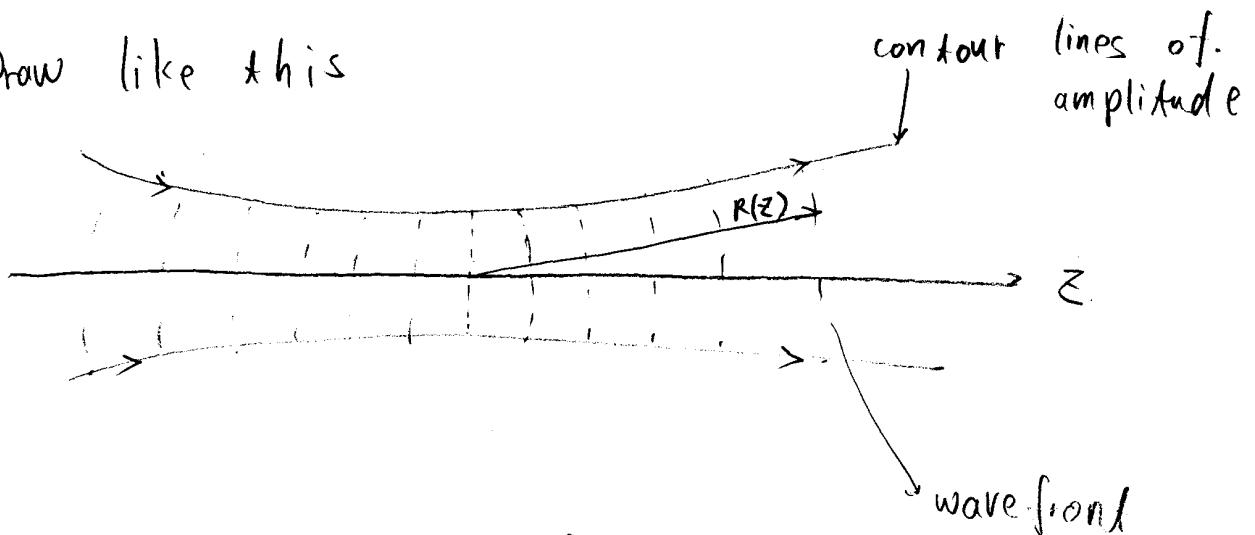
$e^{-ik\theta \frac{p}{zR(z)}} e^{-ikz}$  give phase

↳ Look like paraxial part of a spherical wave.

spherical wave  $\sim e^{-ikr}$   $r = \sqrt{p^2 + z^2}$   
 $= z + \frac{p^2}{2z}$

$R(z)$  = distance to center of curvature.

Draw like this



At focus  $R = \frac{z^2 + z_0^2}{z} = \infty$

wave front is flat

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Far from focus,  $z \gg z_0$   $R \rightarrow z$ .

Wave front centered at focus

So near focus, laser beam looks a plane wave, far from focus.  $\sim$  spherical.

wave centered at focus

Note Laser beam is completely specified.

by two parameters.  $w(z)$  and  $R(z)$

or by complex curvature  $q$

Give  $w$ ,  $R$ , we have

$$\frac{1}{q} = \frac{1}{R} - i \frac{1}{\pi w^2}, \text{ so } q \text{ is determined.}$$

Interesting properties:

$$q(z) = z + i z_0$$

describe beam in terms of  
focus location ( $z$ )  
& beam waist ( $z_0$  or  $w_0$ )

while  $\frac{1}{q(z)} = \frac{1}{R(z)} - i \frac{\lambda}{\pi W^2}$  describes beam in terms of local properties

local width,  $W$   
local radius of curvature  $R$ .

One more contribution to phase

$$u(\vec{r}) = \frac{A_1}{q} e^{-\frac{r^2}{W(z)^2}} e^{-ik(z + \frac{r^2}{2R})}$$

$\frac{1}{q}$  in front has phase too:

$$\frac{1}{q} = \frac{1}{z + iz_0} = \frac{z - iz_0}{z^2 + z_0^2}$$

$$\text{by convention} = -i \frac{z_0 + iz}{z^2 + z_0^2}$$

$$= \frac{1}{\sqrt{z^2 + z_0^2}} e^{-i \frac{\pi}{2}} e^{i\beta}$$

$$\text{where } \tan \beta = \frac{z_0}{z}$$



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$$\text{or } \frac{1}{f} = \frac{e^{-i\frac{\pi}{2}}}{W_0} \frac{W_0}{W(z)} e^{i\beta(z)}$$

Let us write.

$$u(\vec{r}) = A_0 \frac{W_0}{W(z)} e^{-\frac{P^2}{W(z)z}} e^{-ik(z + \frac{P^2}{k(z)})} e^{i\beta(z)}$$

$$A_0 = A_1 \frac{e^{-i\pi/2}}{W_0}$$

$\beta$  called Guoy phase.

varies from  $-\frac{\pi}{2}$  to  $\frac{\pi}{2}$ .