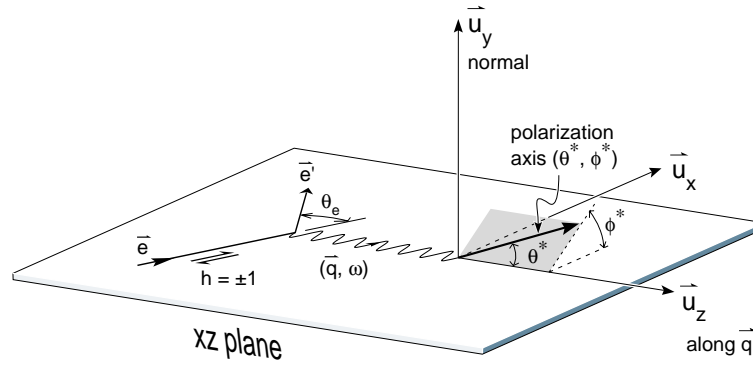


Preliminary Results from the 1998 Run of E93-026

Donal Day

December 8, 2000

Experimental technique for $\vec{D}(\vec{e}, e'n)p$



Cross section:

$$\sigma = \Sigma + h\Delta; \quad h = \pm 1$$

Asymmetry

$$A = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} = \frac{\Delta}{\Sigma}$$

Σ is the elastic unpolarized $e - n$ cross section

$$\Sigma = \sigma_m \frac{E'}{E_0} \left(\frac{(G_E^n)^2 + \tau (G_M^n)^2}{(1 + \tau)} + 2\tau (G_M^n)^2 \tan^2 \left(\frac{\Theta_e}{2} \right) \right)$$

Δ depends on direction of target polarization : Θ^*, Φ^*

$$\Delta = -2\sigma_m \frac{E'}{E_0} \sqrt{\frac{\tau}{1 + \tau}} \tan \left(\frac{\Theta_e}{2} \right) \left\{ \sqrt{\tau \left(1 + (1 + \tau) \tan^2 \left(\frac{\Theta_e}{2} \right) \right)} \cos \Theta^* (G_M^n)^2 + \sin \Theta^* \cos \Phi^* G_E^n G_M^n \right\}$$

Neutron Asymmetry for $\Theta^* = \pi/2$ and $\phi^* = 0$

$$A_n = \frac{-2\sqrt{\tau(1+\tau)} \tan(\Theta_e/2) G_E^n G_M^n}{(G_E^n)^2 + \tau(1+2(1+\tau)\tan^2(\Theta_e/2))(G_M^n)^2}$$

Deuteron asymmetry $A_{ed}^V = \gamma A_n$

$\gamma = 0.92$: correction factor for the D-state of the deuteron

Experimental asymmetries:

$$\varepsilon = \frac{N^{\uparrow} - N^{\downarrow}}{N^{\uparrow} + N^{\downarrow}} = P_B \cdot P_T \cdot f \cdot A_{ed}^V$$

If only life were so simple!

f accounts for the materials other than polarized deuterons in the target.

- The deuteron is, after all, a nucleus and the reaction mechanism is not completely described by PWIA. We must consider the dependence on the reaction mechanism, and the NN potential responsible for the binding.
- Electron interacts with a quasifree nucleon, not a free one, and is ejected over a range of angles which depend its initial momentum and the detector acceptances.
- The experimental asymmetry is distributed over a range in E' , θ_e , ϕ^* and θ^* , $\theta_{nq} \mapsto \theta_{np}^{\text{cm}}$.

Solution: Arenhövel's calculations on a grid over E' , θ_e , ϕ^* and θ^* , combined with a monte carlo that averages the theoretical asymmetries over the detector acceptances.

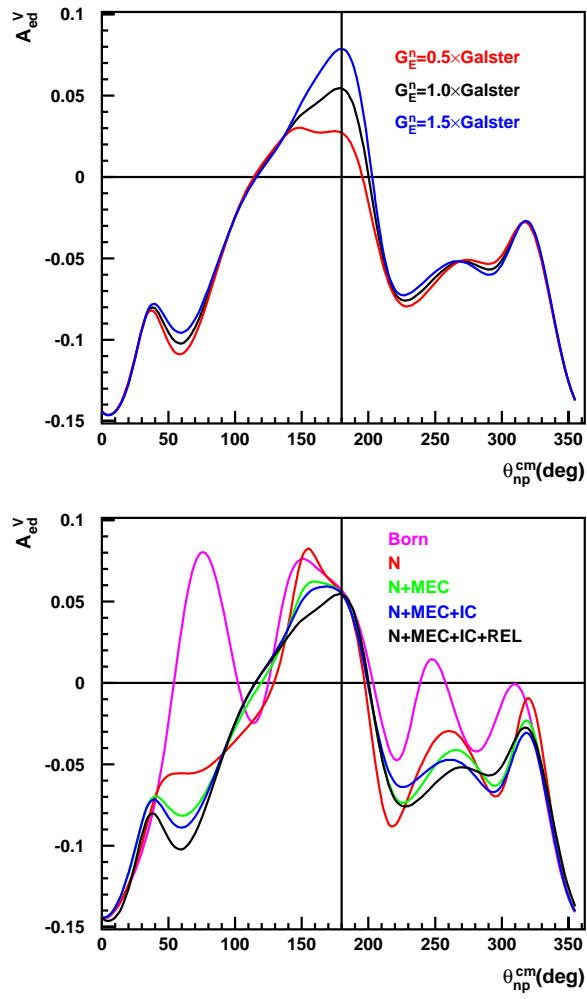


Figure 1: $E=2.721$ GeV, $E'=2.460$ GeV, $\Theta_e=15.8$ deg. Θ_{np}^{cm} : Angle between \vec{q} and relative n-p momentum in cm system. 180 deg corresponds to neutron in direction of \vec{q} .

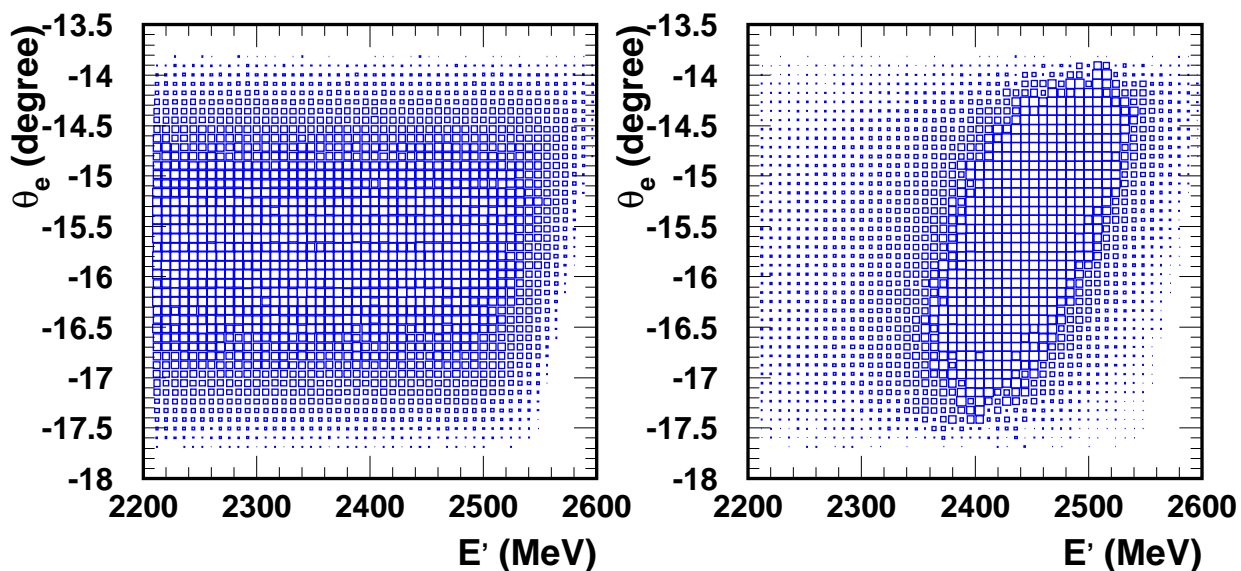


Figure 2: Final electron energy and θ_e distributions from Monte Carlo for quasielastic electron deuteron scattering in E93-026 kinematics. Cross section weighted on right.

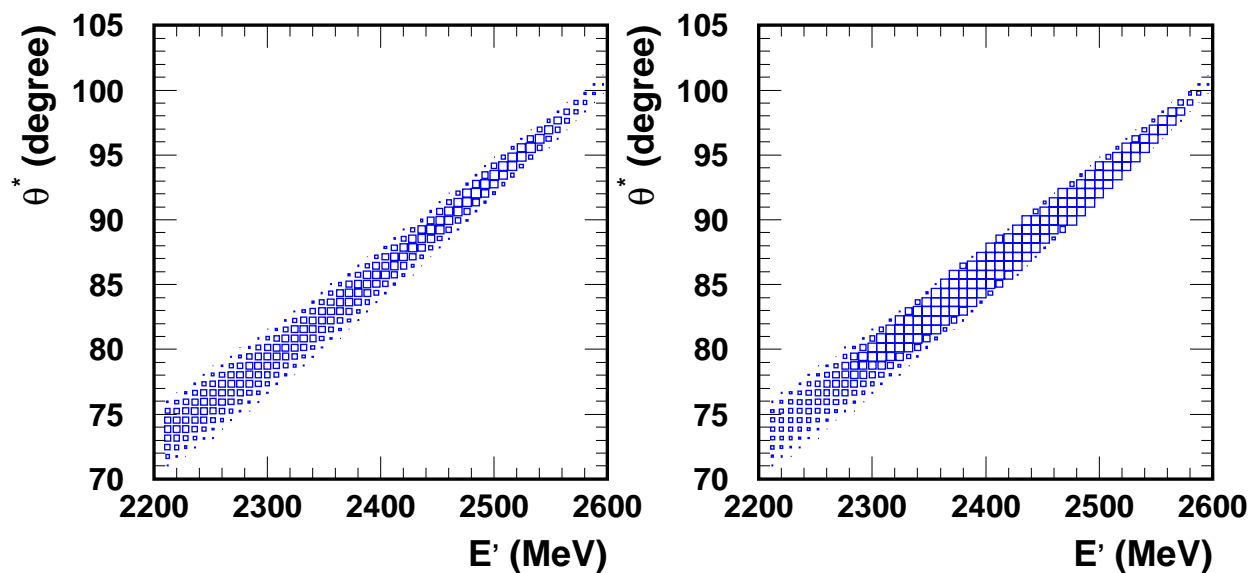


Figure 3: Final electron energy and θ^* distributions from Monte Carlo quasielastic electron deuteron scattering in E93-026 kinematics. Cross section weighted on right.

Download the files

Aved_tht_ep.pdf and Aved_thtcm_tht_e.pdf

from the same directory. The first is a plot of A_v^{ed} against θ_{cm}^{np} and e' at a fixed value of θ_e . The second is A_v^{ed} against θ_{cm}^{np} and θ_e at a fixed value of e' .

f accounts for the materials other than polarized deuterons in the target.

$$f = \frac{n_D \sigma_D}{(1 - \eta_p)n_D \sigma_D + \eta_p n_D \sigma_p + (1 - \eta_N)n_N \sigma_{15} + \eta_N n_N \sigma_{14} + \sum_i n_i \sigma_i}$$

Solution: Monte Carlo and experimental data .

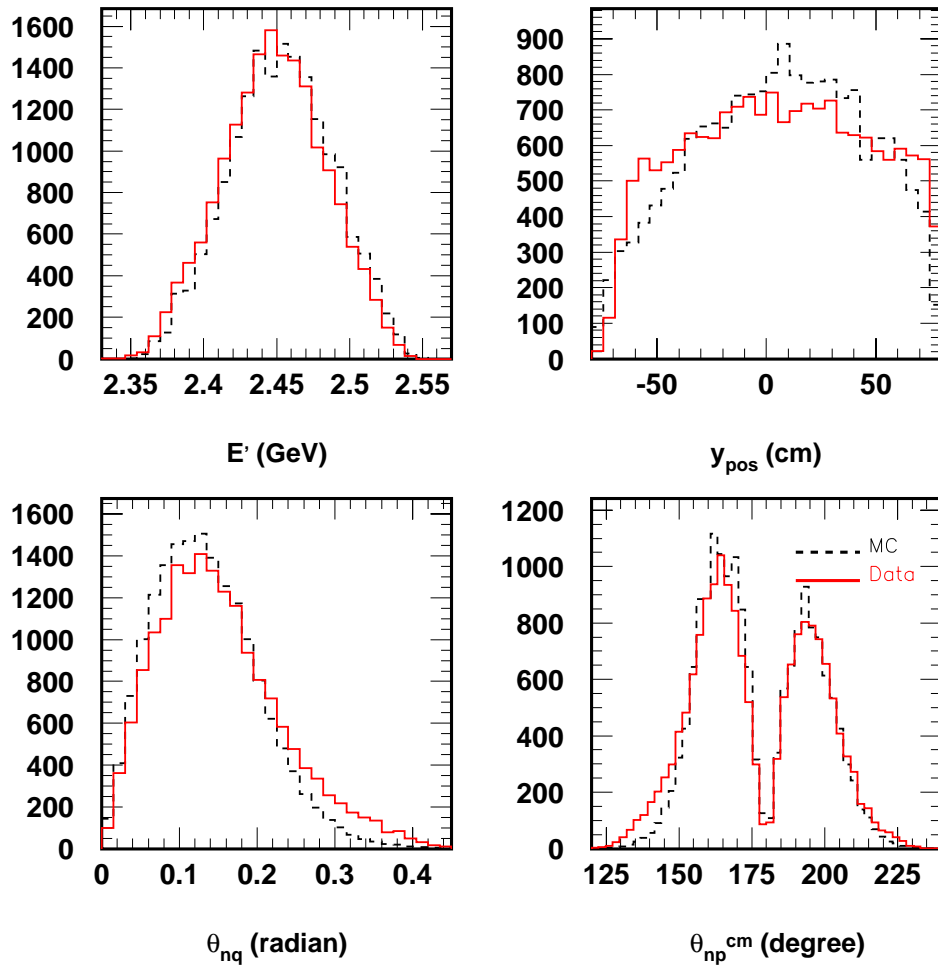


Figure 4: Comparison between data and simulation for (e, e', n) from carbon. Target field, radiative effects included and W cut applied.

The polarized beam and target electron deuteron cross section can be expressed in the form:

$$S(h, P_1^d, P_2^d) = S_0 \left[1 + hA_e + P_1^d A_d^V + P_2^d A_d^T + h(P_1^d A_{ed}^V + P_2^d A_{ed}^T) \right],$$

where $A_e, A_d^V, A_d^T, A_{ed}^V, A_{ed}^T$ are electron asymmetry, vector and tensor target asymmetries, and electron-deuteron vector and tensor asymmetries.

$$A_e = \frac{1}{2hS_0} [S(h, 0, 0) - S(-h, 0, 0)],$$

$$A_d^V = \frac{1}{2P_1^d S_0} [S(0, P_1^d, P_2^d) - S(0, -P_1^d, P_2^d)],$$

$$A_d^T = \frac{1}{2P_2^d S_0} [S(0, P_1^d, P_2^d) - S(0, -P_1^d, P_2^d) - 2S_0],$$

$$A_{ed}^{V,T} = \frac{1}{4hP_{1/2}^d S_0} \left\{ [S(h, P_1^d, P_2^d) - S(-h, P_1^d, P_2^d)] \mp [S(h, -P_1^d, P_2^d) - S(-h, -P_1^d, P_2^d)] \right\}.$$

Note that S_0 is in the denominator and we never measure S_0 .

$$S(+h) + S(-h) \neq S_0$$

.

$$\varepsilon = \frac{(L_+ - R_+) - (L_- - R_-)}{(L_+ + R_+) + (L_- + R_-)}.$$

$$\begin{aligned} L_+ &= \Phi_+ n_D S(h, P_1^d, P_2^d) \\ L_- &= \Phi_- n_D S(h, -P_1^d, P_2^d) \\ R_+ &= \Phi_+ n_D S(-h, P_1^d, P_2^d) \\ R_- &= \Phi_- n_D S(-h, -P_1^d, P_2^d) \end{aligned}$$

$$\varepsilon = \frac{h [(1 - \beta)A_e + (1 + \alpha\beta)P_1^d A_{ed}^V + (1 - \beta\gamma)P_2^d A_{ed}^T]}{(1 + \beta) + (1 - \alpha\beta)P_1^d A_d^V + (1 + \beta\gamma)P_2^d A_d^T},$$

$$\alpha = -\frac{P_{1-}}{P_{1+}}$$

$$\beta = \frac{\Phi_-}{\Phi_+}$$

$$\gamma = \frac{P_2(P_{1-})}{P_2(P_{1+})}$$

$$P_2^d = 2 - [4 - 3(P_1^d)^2]^{1/2}$$

$$A_{ed}^V = \frac{1}{h(1 + \alpha\beta)P_1^d} \left\{ \varepsilon [(1 + \beta) + (1 - \alpha\beta)P_1^d A_d^V + (1 + \beta\gamma)P_2^d A_d^T] - h [(1 - \beta)A_e + (1 - \beta\gamma)P_2^d A_{ed}^T] \right\}.$$

With full ϕ acceptance A_e , A_d^V and A_{ed}^T have zero contributions.

A_d^T remains. For $P_1^d = 20\%$, $P_2^d \simeq 3\%$

with $A_d^T \approx 10^{-2} A_d^T$ can be ignored

$$A_{ed}^V = \frac{\varepsilon(1 + \beta)}{h(1 + \alpha\beta)P_1^d}.$$

With the dilution factor

$$A_{ed}^V = \frac{\varepsilon(1 + \beta)}{h(1 + \alpha\beta)P_1^d f}.$$

ANALYSIS

- **Event Reconstruction**

- HMS standard reconstruction + target magnet + raster
- NDET tracking: Particle Identification

⇒ **variables:** $w, Y_{pos}, cointime, \theta_{nq}$

- **Event Selection and Normalization**

→ Normalized yields: N^\uparrow, N^\downarrow

- **Experimental asymmetries:**

$$\varepsilon = \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} = P_B \cdot P_T \cdot f \cdot A_{ed}^V$$

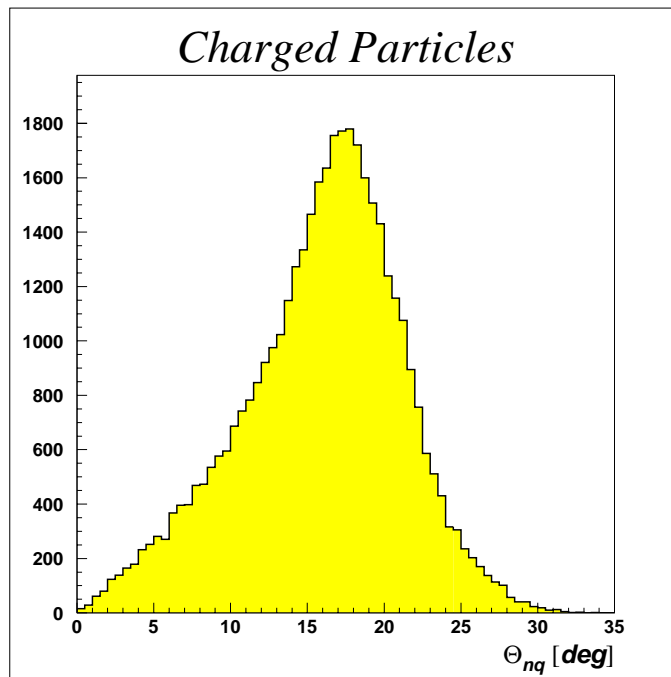
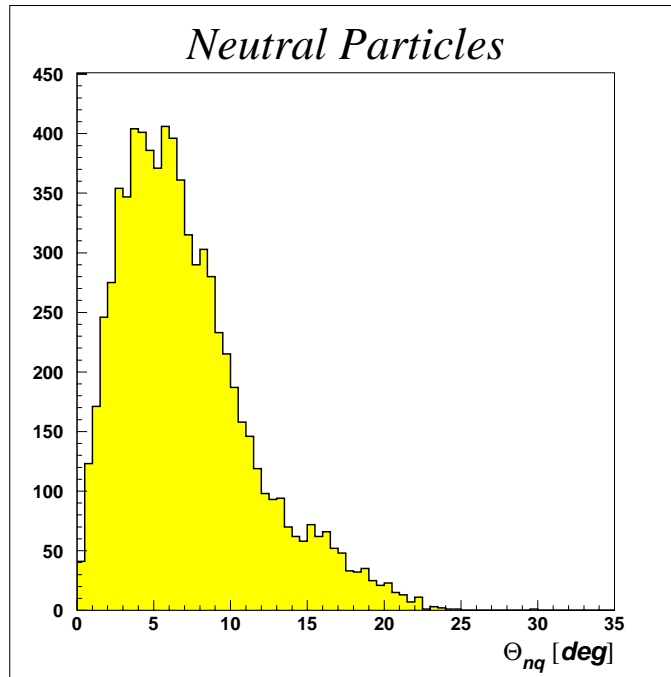
- Beam Polarization P_B from moller polarimetry, $\simeq 77\%$
- Target polarization P_T : average: 22%
- Dilution factor $f = \frac{N_{polarized}}{N}$ from ^{12}C data and MC

- **Radiative corrections, Accidental background subtraction**

⇒ **measured deuteron asymmetry** A_{ed}^V

- **Theoretical calculations of A_{ed}^V (for different G_E^n models) from Arenhövel for a kinematical grid**
- **Average theoretical A_{ed}^V over experimental acceptance using MC and compare with measured A_{ed}^V**

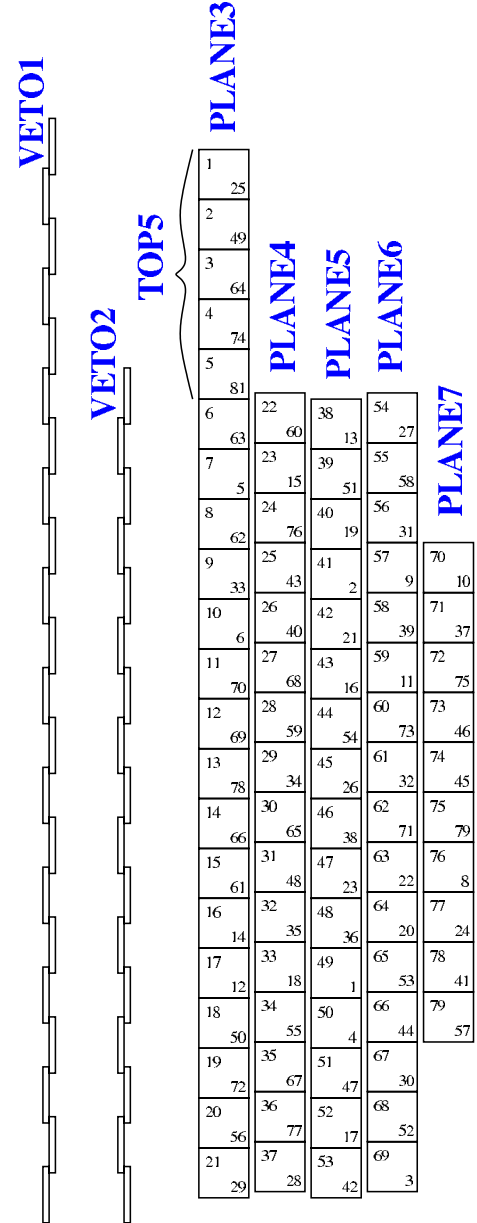
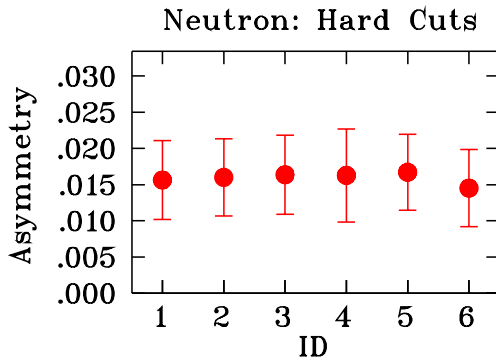
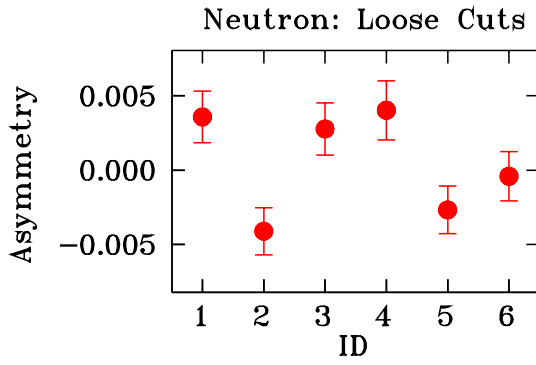
⇒ G_E^n



NEUTRON IDENTIFICATION

Definitions:

- ID1 = Analyzer Neutron
- ID2 = (!VETO1) && (!VETO2)
- ID3 = ID2 && (!TOP5)
- ID4 = ID2 && (!PLANE3)
- ID5 = (!VETO1) && (!TOP5)
- ID6 = (!VETO2) && (!TOP5)



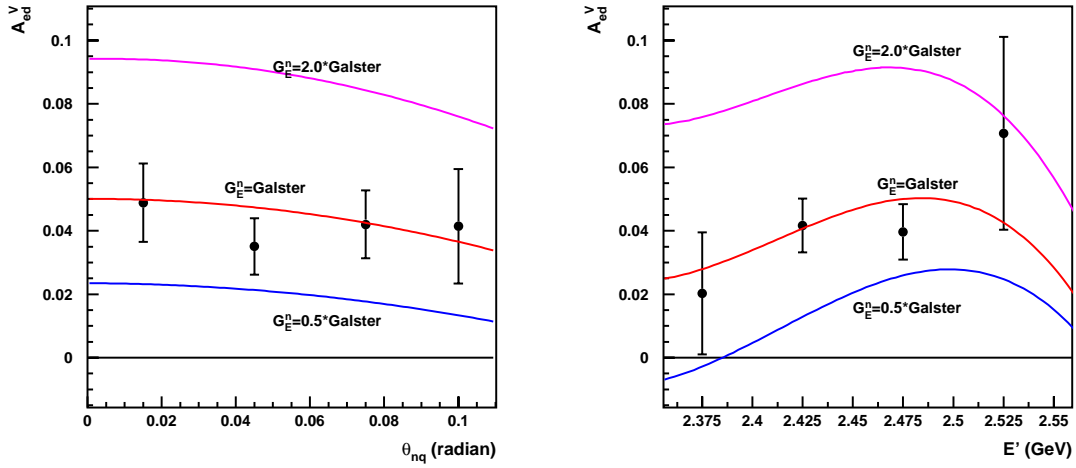


Figure 5: Right figure: E' : Energy of the scattered electron. Left figure: θ_{nq} : Angle between 3-momentum transfer \vec{q} and hadron track

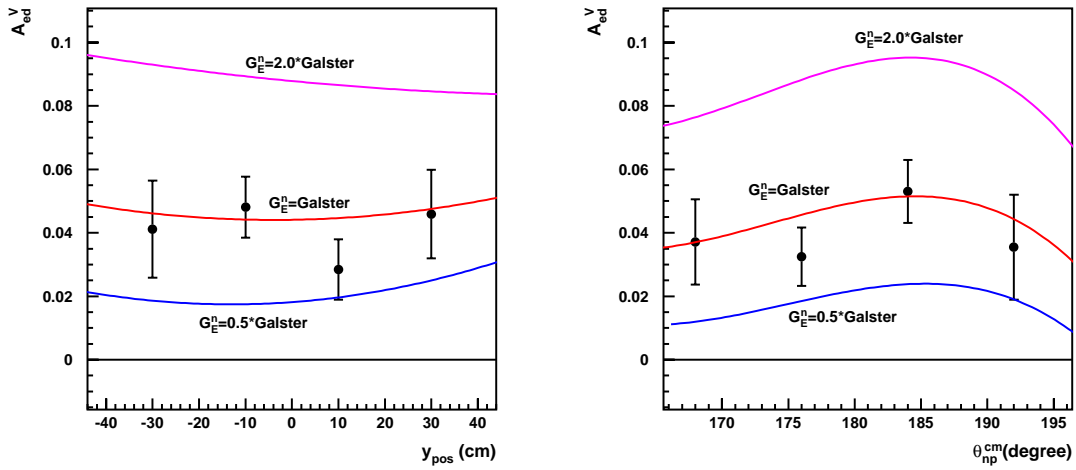
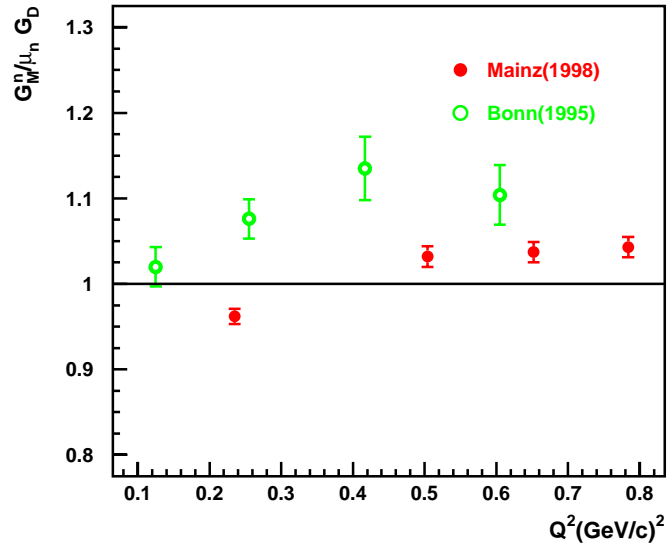


Figure 6: Left Figure: Y_{pos} : Horizontal coordinate of the hadron track in the neutron detector. Right Figure: θ_{np}^{cm} : Angle between 3-momentum transfer \vec{q} and relative n-p momentum in the cm system (180 deg corresponds to neutron in direction of $vec{q}$)

Impact of G_M^n

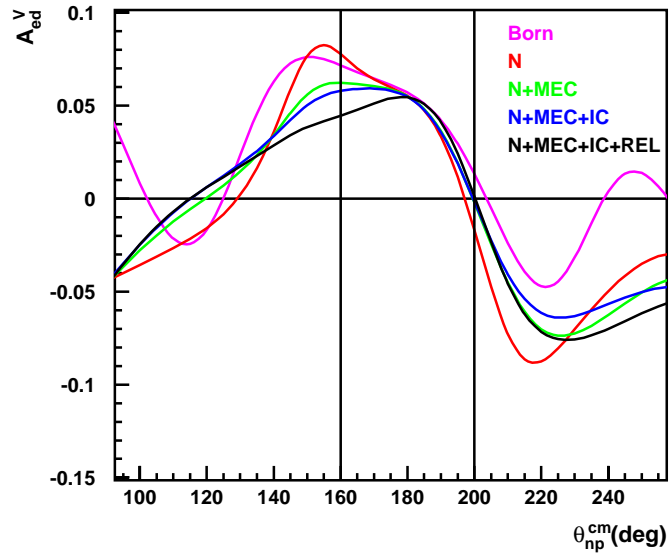


Difference between $G_M^n = \mu_n G_D$ and

G_M^n Mainz: **3%**

G_M^n Bonn: **10%**

Reaction Mechanism Dependence



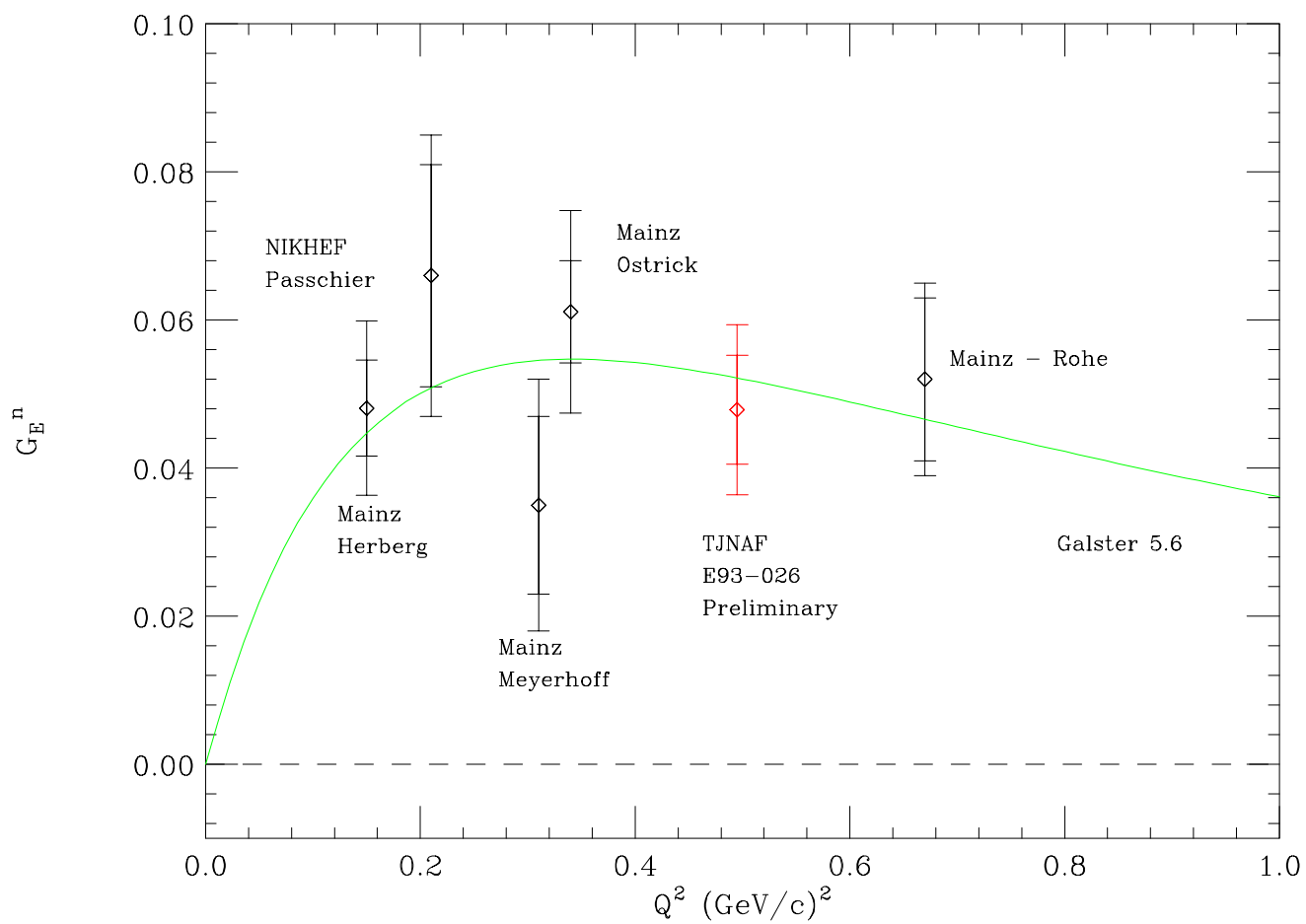
Experimental acceptance: $160 \text{ deg} < \theta_{np}^{cm} < 180 \text{ deg}$

Difference between full Calculation and

Born: 14%

Born + REL: $\simeq 8\%$

Polarized Experiments



RESULTS

Preliminary result for $Q^2 = 0.5 (GeV/c)^2$:

$$G_E^n = 0.04788 \pm 0.00737_{stat} \pm 0.00378_{syst}$$

Systematics:

Source	Contribution	status
G_M^n	3.3%	✓
Beam Polarization	1.1%	✓
Target Polarization	5.3%	✓
Dilution factor	5%	✓
Radiative Corrections	1.0%	✓
A_{ed}^V averaging	0.4%	✓
Kinematic Uncertainties	1.9%	✓
NDet position offset	1.4%	✓
Total $(\Delta G_E^n / G_E^n)_{syst}$	7.9%	✓